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## Electrochemical Time of Flight for Rapid and Direct Measurement of Diffusion Coefficients

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# Electrochemical Time of Flight for Rapid and Direct Measurement of Diffusion Coefficients

A dissertation submitted in fulfillment  
of the requirements for the degree of  
Doctor of Philosophy in Chemistry

by

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Valparaiso University  
Bachelor of Science in Chemistry, 2012

May 2017  
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This dissertation is approved for recommendation to the Graduate Council.

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## Abstract

The determination of diffusion coefficients is of fundamental importance to the understanding of electrochemistry and sensors. Developing a method by which diffusion coefficients of Red/ox active analytes can be determined quickly and elegantly, would be a great advancement over presently accepted methods. This dissertation reports the reviving electrochemical time of flight (ETOF), and developing a method that allows for empirical determination of diffusion coefficients from a single measurement. ETOF is a generate and detect experiment where the time an electrochemically generated species takes to transit a known distance is measured and related to the diffusion coefficient of the species. The determined diffusion coefficient of ferricyanide,  $7.3(\pm 0.7) \times 10^{-6} \text{ cm}^2/\text{s}$ , was within the 95% confidence interval of the literature value, using the traditional ETOF data treatment. In this dissertation a new treatment of the data, the Moldenhauer treatment, where a diffusional calibration curve is constructed using multiple species of known diffusion coefficient and measuring their transit times at a set distance. The calibration curve constructed in aqueous solutions found the diffusion coefficient of ruthenium(II) hexamine to be within the 95% confidence interval of what has been reported in the literature. The same calibration was also used to determine diffusion coefficients of aqueous probe molecules in a more viscous solution of 20% v/v ethylene glycol and water. Computational modeling was used to further optimize generator pulse widths to allow for a greater linear range of determineable diffusion coefficients. It was shown that an empirically determined aqueous calibration can be used to determine diffusion coefficients in organic solvents. The diffusion coefficient of ferrocene was determined to be  $2.4(\pm 0.1) \times 10^{-5} \text{ cm}^2/\text{s}$  after modeling directed optimum generator pulse widths. In addition diffusion coefficients were determined for tetrabutylammonium dioxovanadium(V) dipicolinate ( $3.9(\pm 0.2) \times 10^{-6} \text{ cm}^2/\text{s}$ ), and ruthenium (II) bisbipyridine dichloride ( $9.3(\pm 0.4) \times 10^{-6} \text{ cm}^2/\text{s}$ ), which do not have published diffusion coefficients presently. In the future, this same method could be used to determine diffusion coefficients in membranes and complex solvents such as ionic liquids.

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## **Dedication**

I would like to dedicate this dissertation to several people. Firstly, my significant other Cassandra Gronenedyke, for having had to deal with me when research is going well and when research was going poorly. Secondly, to my laid back boss, Dr. David Paul, and his catchphrase “Dumb looks are Free”. And finally to my parents who at least only minimally asked any of those questions that grad-students dread hearing.

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## **List of Published Papers**

Chapter 1: " Rapid and Direct Determination of Diffusion Coefficients Using Microelectrode Arrays "  
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## Introduction

Diffusion is both an important property to electrochemistry and a property commonly determined using electrochemistry. Diffusion of molecules through a solvent is important to measure when trying to elucidate electrochemical mechanisms (1-9) for, design of better batteries and other energy applications(10-22), electrochemical sensors(23-29), optimization of industrial processes(30), and for calculating conductive properties of electrolytes and solutions (31-43). All of these processes have steps that are limited by the mass transport of chemical species, from the most fundamental research in electrochemical measurements to more applied areas of research such as energy storage and conversion or medical sensors(44) the values of diffusion coefficients are of great importance . Diffusion coefficients are everywhere because mass transport is everywhere and measuring them would add to our understanding of the world around us.

In order to better understand the importance of research into methods of determining diffusion coefficients, it is best to discuss an application for which diffusion coefficients need to be measured, that area being a growing area of electrochemistry: room temperature ionic liquids (RTILs). These liquid organic molecules are solvents that also serve as a supporting electrolyte for electrochemical experiments. RTILs have an advantage over other liquid salts in that they can be studied in a lab setting without special high-temperature equipment. However, RTILs are an area of great interest in electrochemistry right now, and because of their solvation properties and wide potential windows, they provide a new area for electrochemical study, and few diffusion coefficients for molecules in RTIL's have been published. Diffusional properties of species in these solvents is important for determining how well they can be used for electrocatalyzed synthesis(40), such as electrocatalyzed carboxylation(3), and as a solvent for other types of electrochemical measurements for physical chemistry and mechanistic studies(8, 41, 42). RTILs can serve as the best options for electrolyte solutions in batteries(20) where output currents are limited by diffusion between the anode and the cathode(38). Faster diffusion translates to less impedance to current across the separator(12). Measuring diffusion coefficients in other solvents such as high temperature ionic liquids that could be used in large storage batteries for power plants and bulk power generation(4, 31, 45), as well as for their use in fuel cells would also be important(14). Knowledge of diffusion coefficients in RTILs also contributes to the understanding of

mechanisms behind various photocells such as dye-sensitized solar cells(19, 43). Ionic liquids also have very large and adjustable potential windows and that leads them to be exciting as new solvents for electrochemical systems(36). Diffusion coefficients shift based on the purity of ionic liquids(37), for instance the water content in the ionic liquid(39), so the possibility of determining purity via diffusion coefficients exists.

For all of these reasons there needs to be a way to quickly and easily determine diffusion coefficients. Electrochemical Time of Flight in the literature appears to be a quick and effective method of determining diffusion coefficients. However, no one, yet, has used ETOF for anything more than a method to provide empirical backing for computational diffusion models. This dissertation shows that ETOF can be used for the determination of diffusion coefficients of Red/Ox species and diffusion coefficients in unlike solvents. This broadens the applicability of a method that has been mostly ignored since its inception. ETOF can be used to determine diffusion coefficients with a simple measurement and in less time than traditional methods. If broad application could be shown, as it has in this dissertation, researchers would be able to use it as a common method for determination of diffusion coefficients.

### **A Comparison of Methods for determining diffusion coefficients**

One of the most common ways to evaluate the diffusion of Red/Ox active molecules in a bulk solution is to use rotating disk electrodes (RDE) and the Levich equation, Equation 1 relating the diffusion of a species to its limiting current(46).

$$i_{lev} = 0.620nFAD^{\frac{2}{3}}\omega^{\frac{1}{2}}\nu^{\frac{-1}{6}}C_s \quad \text{Equation 1}$$

Where the limiting current,  $i_{lev}$ , is proportional to ; $n$  the number of electrons transferred,  $F$  Faraday's constant,  $D$  the diffusion coefficient,  $\nu$  the kinematic viscosity of the supporting electrolyte,  $\omega$  the rotation rate, and  $C_s$  the analyte concentration in bulk(47). This method requires one to have prior knowledge of the concentration of the analyte in question, the electrode area, and the number of electrons transferred, the viscosity of the system, and electron transfer kinetics. So this method works well only for fully elucidated systems such as any number of common probe molecules, i.e. ferricyanide, ferrocyanide, etc.. There are several limitations to determining the diffusion coefficient of an analyte using the Levich

equation. There is an issue with very viscous solvents such as some ionic liquids where RDE techniques fail because true diffusion controlled limiting currents are difficult to achieve (48). However the Levich Equation is still the standard method in conventional solvents, even though that concentration is the only known and controlled factor (49). Sometimes the number of electrons transferred is an unknown quantity, so an alternative method to determining diffusion coefficients is needed. One such alternative is ETOF, which is the only technique capable of determining diffusion coefficients without foreknowledge of the concentration of the analyte, the electroactive area, the number of electrons transferred, or the viscosity of the solvent(50), and in small volumes.

Royce Murray and his group were the first to use the ETOF experiment in the late 1980s(51). They were interested in the apparent diffusion of electrons across conducting polymers. They sandwiched a conducting polymer between two electrodes and monitored a pulse of charge as it travelled between them. This they related to the Einstein equation, Equation 2, where  $d$  is the distance traveled (space between the electrodes),  $\theta$  is a numerical constant that is related to the geometry of the electrodes,  $D_e$  is the diffusion coefficient of an electron through the polymer, and  $t_{max}$  is the time at which there was a maximum current at the detecting electrode.

$$\frac{D_e t_{max}}{d^2} = \theta \quad \text{Equation 2}$$

ETOF then was applied to determining of the diffusion coefficient of an analyte in bulk solution. In 1990 Stuart Licht(52) published a paper which laid out a method for using individually addressable microelectrode array and ETOF to determine the diffusion coefficient of a molecule without first knowing its concentration, simply by knowing the distance and the time that the electrochemically active species took travel between the generator and the detector electrode. From his work he determined that the diffusion coefficient ( $D$ ) was inversely proportional to the time of maximum collection and that the time of maximum collection ( $t_{max}$ ) was directly proportional to the distance ( $d$ ) squared as related by an empirically determined proportionality constant, Equation 3 (52).

$$t_{max} = 0.22 \frac{d^2}{D} \quad \text{Equation 3}$$

It is important to note that the time of maximum collection is measured from the middle of the generation pulse to the peak on the collector current. The distance,  $d$ , is measured from the center of the generator electrode to the edge of the collection electrode. He also modeled this equation using a random walk simulation and arrived at an equation that appears to be in close agreement with the above,  $t_{max} = 0.20 \frac{d^2}{D}$ , which he stated had an experimental uncertainty of 10%(52).

Christian Amatore(50) continued to look at the above method as a way of determining diffusion coefficients of analyte molecules in a bulk solution explaining why pulsed generation is better than continuous generation. This is because it is easier to determine the time of maximum collection current than the time as which a steady state current is achieved. In addition, when the generator is polarized for longer periods of time there is a greater chance of molecules from the detector diffusing back to the generator(50). This diffusion from the collector back to the generator (sometimes called diffusion layer over-lap or cross-talk) lengthens the amount of time it takes to achieve the limiting current, and determining the  $t_{max}$  is difficult. Longer times also require very stable redox species to serve as the analyte(50). Amatore also reconstructed Licht's equation, Equation 4, where  $d$  is still the distance,  $K$  is the geometry constant for the electrode system that must be determined empirically for each electrode system,  $D$  is the diffusion coefficient of the detected species, and  $t_{max}$  is the time of maximum collection(50).

$$d = K\sqrt{Dt_{max}} \quad \text{Equation 4}$$

Amatore correctly determined the diffusion coefficients for the ferri-ferrocyanide system in bulk solution by ETOF using the above equation after finding  $K$  for the electrode array.

Despite there being some skepticism about generate and detect experiments working in low electrolyte concentrations(53), with slight experimental modifications, electrochemical time of flight can determine a variety of other interfacial parameters. Using galvanic generation, Slowinska was able to determine the capacitance of solid films on potentiometric sensors. The sensing method at the detector is potentiometric, the method being called potentiometric electrochemical time of flight or P-ETOF(54-56). The same group that did the P-ETOF experiments also applied electrochemical time of flight to viscous

solutions and showed the ability to determine diffusion coefficients of analyte molecules as they change with increasing solution viscosity(57). Others have used it for the determination of diffusion coefficients in glucose solutions and gels(58) and still another group has shown that ETOF can actually be used for the determination of the diffusion coefficient through a solid substrate(59).

However there are other methods of determining diffusion coefficients that include both electrochemical and non-electrochemical means. Other electrochemical means of determining diffusion coefficients include, chronoamperometry (25, 35) and scanning electrochemical microscopy (SECM)(34). The Cottrell equation (Equation 5) relates current,  $I(t)$ , to the time,  $t$ , and concentration of the analyte,  $C_0$ , in solution.

$$I(t) = \frac{nFAD_0^{\frac{1}{2}}C_0^*}{\pi^{\frac{1}{2}}t^{\frac{1}{2}}} \quad \text{Equation 5}$$

Unfortunately, it has many of the same limitations as the Levich equation, requiring knowledge of the electroactive area of the electrode,  $A$ , and the number of electrons transferred,  $n$ . The equation can be used in both voltammetric and amperometric experiments but more commonly amperometric measurements of diffusion coefficients are performed.

Another method similar to ETOF that can be used to determine diffusion coefficients is scanning electrochemical microscopy (SECM). This uses an ultramicroelectrode (UME) electrode held a short distance vertically from a surface, or substrate, and the UME, or tip, are used as a detector and a generator electrode respectively and biased so that opposing reactions are happening at each electrode(60-65). SECM is typically used for surface studies of electrode arrays and thin films, both making use of the fact that the tip can be moved and determine spatial data about the resistive or conductive nature of the substrate. It can also be used for kinetics studies and the determination of diffusion coefficients(60, 62, 65). Determining the diffusion coefficient ratio of the couple in the solution from the limiting currents at the detector and from the feedback to the generator is based on knowing the current at time equals infinity, the electroactive area of the electrode, and the starting concentration of one the analytes (61). This experiment returns to a methodology that requires knowledge of the electroactive area of the electrode and the initial concentration of one of the species, however it does not require knowledge of how far electrodes are separated.

There are still other non-electrochemical options, such as pulsed field gradient NMR, which one can use to determine the diffusion coefficient of a species. In these experiments diffusion of the ions such as lithium, that are detectable by NMR, is measured by the decay of echo intensity precession of the nuclei on the atoms to the RF excitation(12). This method is commonly used for lithium ions to examine gel electrolytes or other materials for use in batteries(11-13). This makes it a popular method for using with viscous organic solvents like the ionic liquids previously mentioned that are liquid organic anions and cations, because they are common for use as solvents in batteries and because of their viscosity it becomes difficult to use for hydrodynamic experiments(11). This method is also primarily used to determine the ionic conductivities of organic solvents because it is easy to measure the diffusion coefficient of atomic ions(32). This method can also be used for organic molecules diffusing in ionic liquids(33), the solubility and diffusion of gas molecules such as CO<sub>2</sub> in ionic liquids(30), and organic molecules that are not electrochemically active(66).

The power of ETOF is that as opposed to hydrodynamic methods it can be used over a wide range of viscosities and use much smaller volumes. Meaning that the Moldenhauer treatment, constructing diffusional calibration curves, could be used for determination of diffusion coefficients in highly viscous ionic liquids, but unfortunately a limited amount of information on diffusion coefficients for probe molecules in ionic liquids and liquids of increased viscosity exists. Such a calibration curve could not be constructed without knowledge of diffusion coefficients in the specific solvent system used. What was revealed in this study was that it is possible to construct a diffusional calibration curve in one solvent system using a species of known diffusion coefficients; then use that same calibration curve to determine the diffusion coefficients of molecules in different solvent systems by simply measuring the TOF and using the calibration curve to determine the diffusion coefficient. This is something that has never been shown in by showing that ETOF can be applied to numerous unlike compounds in unlike solvents using an aqueous calibration curve, a stepping stone for research in other solvent systems.

### **Summary of Presented Work**

ETOF methods as we present it here has never before been used before to determine diffusion coefficients of multiple unlike species. In fact while it has been used empirically in the past it has never

been previously utilized beyond finding a diffusion coefficient of a single molecule or to verify diffusional models of electrochemical experiments. The work here provides the groundwork needed so that one could use this as a primary technique for the determination of diffusion coefficients, and especially to apply this technique to the diffusion of molecules through membranes and for the determination of diffusion coefficients through ionic liquids. Although these applications are important, only diffusion coefficient in bulk solution were determined here.

The first paper of this volume shows that one can make an empirical calibration curve by rearranging the ETOF equation and describes the construction of the Generation Electrode Controller (GEC) that was used to control the voltage pulse on the generation electrode. The GEC is controlled by a National Instrument's LabVIEW program providing for pulsed generation on the second working electrode of a CHI 750series potentiostat, more details on the construction of the device can be found in appendix 1. We were then able to determine diffusion coefficients using an empirical calibration of a microelectrode array using both the traditional treatment of ETOF using as single analyte and our non-traditional data treatment using a series of standards to construct a calibration curve for the determination of diffusion coefficients based on their time of flight. The curves in this paper suggested that there was an intercept that we later determined was unimportant.

In the second paper we approached the question of: What is the range of diffusion coefficients that can be determined by using our alternative data treatment (with any accuracy)? For that we collaborated with Christian Amatore and Catherine Sella to model our experiment to optimize the parameters over a wide range. Here we report for the first time that  $K$ , is dependent on the pulse width at the generator as well as the gap between the generator and the collector.  $K$  is constant only for a range of generation pulse widths and that as long as the generation pulse is tuned correctly, one can determine diffusion coefficients for a wide range of analytes across at least 1 order of magnitude using a single electrode geometry. This paper also shows that diffusion coefficients can be determined using the Moldenhauer data treatment in vastly different solvent systems and viscosities of solutions, which was something that had not previously been reported in the literature.



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## **Rapid and Direct Determination of Diffusion Coefficients Using Microelectrode Arrays**

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**Abstract:** The sensitivity of amperometric sensors is typically set by the rate diffusion of the analyte to the electrode surface, so determining diffusion coefficients in various electrolyte solutions is of fundamental interest. It has been theoretically shown and verified that diffusion coefficients of electrochemically generated analytes can be determined using electrochemical time of flight (ETOF), a method that uses an electrochemical array in which one electrode generates a Red/Ox species, and measures the analyte diffusion times to collecting electrodes of differing distances from a stationary generator. ETOF has the potential to greatly simplify the determination of diffusion coefficients as the analyte concentration, the electroactive area, the solution viscosity, and the electron transfer kinetics can remain unknown. Here we demonstrate a rearrangement of the ETOF experiment in which the electrochemical flight time is measured for a series of different Red/Ox species of known diffusion coefficients at single distance. We show this a valid application of a method that has existed for almost 30 years, by determining diffusion coefficients for ruthenium (II) hexamine, and diffusion coefficients in solutions of increased viscosity.

## Introduction

Diffusion coefficients are important because they set the sensitivity of amperometric sensors and they are a fundamental property both in membrane permeability and in electrochemical measurements. The most common method of determining diffusion coefficients for analytes in bulk solutions or through gels and membranes relies on the rotating disk electrode (RDEs) (1-7) or the rotating ring disk electrode (RRDE) (8). This method determines the diffusion coefficients,  $D$ , from the slope of a Levich plot constructed by measuring limiting currents,  $I_L$ , as a function of square root of the rotation rate,  $w$ , according to the Levich equation (Equation 1).

$$I_L = 0.620nFAD^{\frac{2}{3}}w^{\frac{1}{2}}v^{-\frac{1}{6}}C \quad \text{Equation 1}$$

Accurate values for the area of the electrode,  $A$ , the number of electrons transferred,  $n$ , the concentration of the molecule,  $C$ , and the viscosity of the solution,  $v$ , must also be known in order to effectively determine the diffusion coefficient from the slope of a Levich Plot. The diffusion coefficients of molecules through bulk solution can also be determined quantitatively by wall-jet chronoamperometry(9), or qualitatively by comparing the CV's of different compounds because the shape of the CV is related to the diffusion coefficient of the molecule (10-12). The other primary option for determining diffusion coefficients of a molecule through a membrane coated over an electrode is impedance spectroscopy (13-18), where the diffusion of the molecule through a membrane or polymer is related to the impedance of the polymer or membrane to current flow. As such, the diffusion through the polymer is related directly to the resistance of charge transfer (mobility) through the membrane, which is related to its conductivity and directly correlated to the diffusion coefficient by the Nernst-Einstein equation (Equation 2).

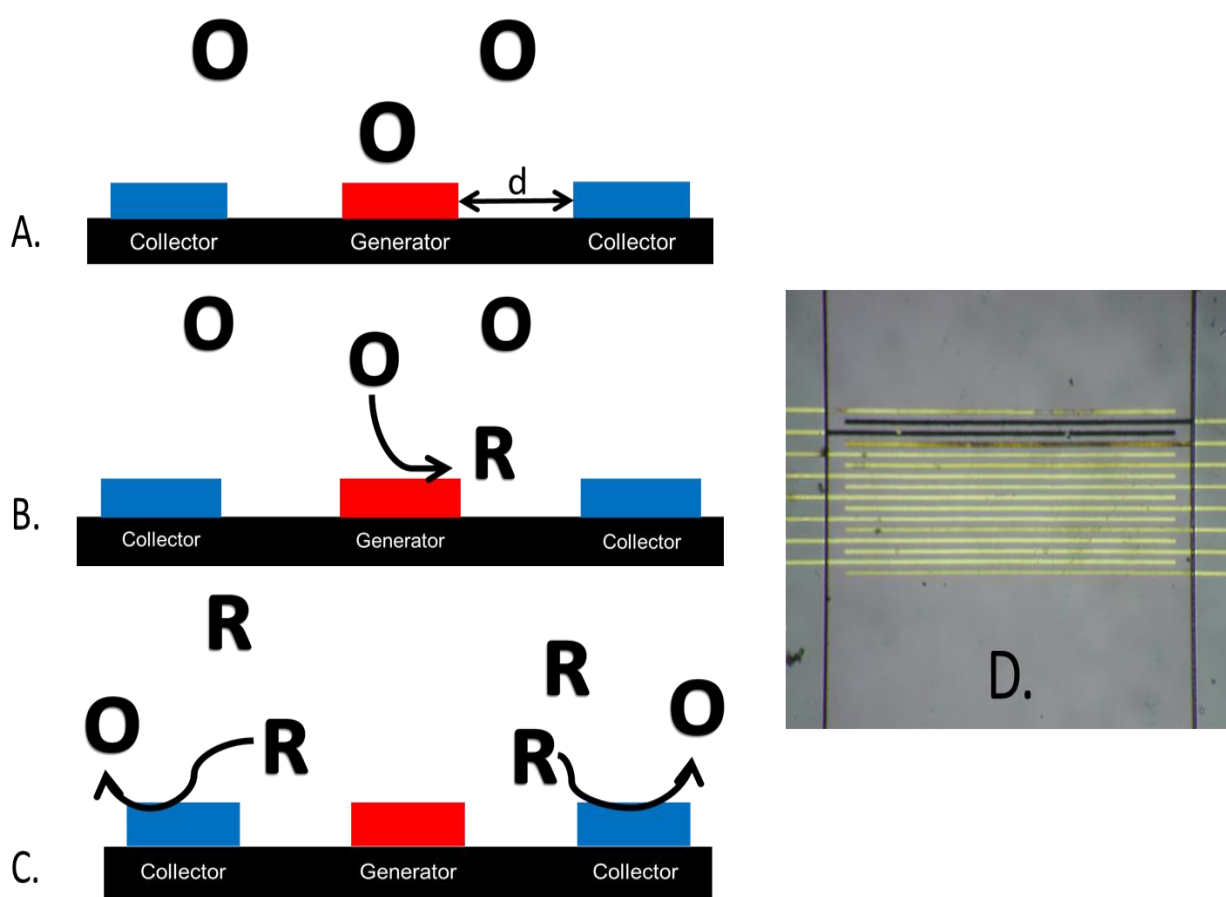
$$D = \frac{\sigma kT}{Cq^2} \quad \text{Equation 2}$$

Conductance,  $\sigma$ , can be determined from the charge transfer resistance and related to the diffusion coefficient for the analyte,  $D$ , if the concentration,  $C$ , the temperature in Kelvin,  $T$ , and the charge on the species,  $q$ , are also known, where  $k$  is the Boltzmann Constant. The above equation is primarily used

when looking at diffusion of ions through solid polymer electrolytes, membranes, and polymer brushes (13-16).

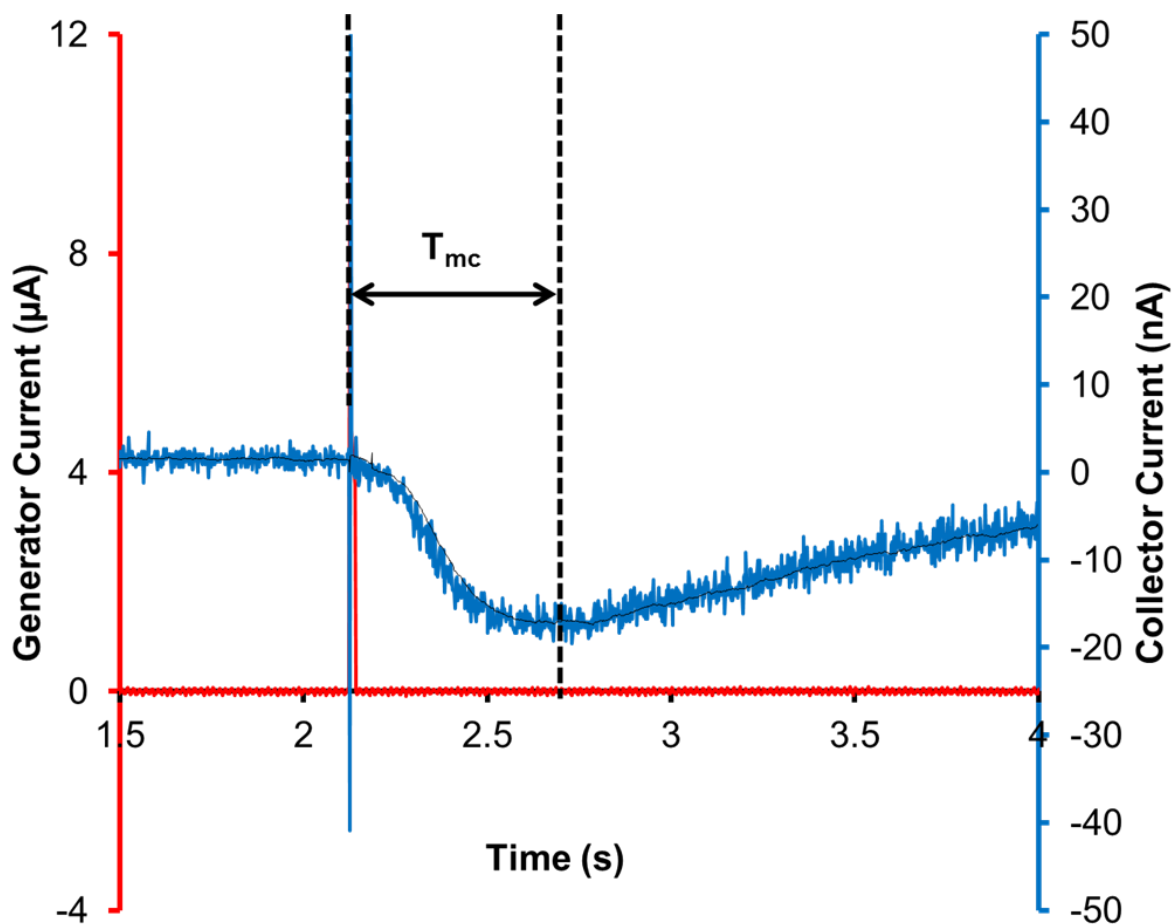
Methods using conductance cells for determining the diffusion coefficients of ions in a solution use capillary flow tubes: one capillary containing 25% more concentrated and one containing 25% less concentrated solution than the bulk. The change in the ratio of the resistances of the solutions in the two flow tubes as the ions diffuse from the capillaries into the solution is measured. These ratios can be modeled by Onsager-Fuoss to determine the diffusion coefficient of an electrolyte diffusing in bulk solution (19-24). There are also non-electrochemical methods for determining diffusion coefficients such as: NMR, field-flow fractionation, and neutron radiography (25). These determinations are very involved, as for each candidate an entirely new experiment must be developed, and numerical values for the ancillary parameters must be known, either taken from the literature or measured. For example, the slope of the Levich line depends on viscosity. Slight changes in the viscosity, such as differences in the composition of the supporting electrolyte, can cause shifts in the diffusion coefficient. In addition, this kind of detailed work has only been done for a few probe molecules that are electrochemically ideal in their behavior. For example, the diffusion coefficients of potassium ferricyanide is known in several buffers, but if/when the buffer is changed, the diffusion coefficient must be re-determined. In addition, none of these molecules are typically the desired analytes for electrochemical sensors, where the sensitivity of the sensor is related to the diffusion of the analyte to the electrode or through the membrane coating the sensor (26). It would therefore be convenient to find a way to determine the diffusion coefficients of Red/Ox analytes in supporting electrolyte solutions.

In this paper we show that there is an existing method in the literature that has been mostly ignored in the thirty years since its development. Electrochemical-time-of-flight (ETOF) is a generate-detect experiment in which an analyte is generated either oxidatively or reductively at one electrode, called the generator; it is detected by re-reducing/re-oxidizing the analyte at a second electrode, called the collector (Figure 1).



**Figure 1:** The Electrochemical Time of Flight experiment (ETOF). A) Oxidized form (O) in solution and generator (red) is at open circuit, the collectors (blue) are polarized to an oxidizing potential. B) The generator is briefly polarized to a reducing potential, converting the O to its reduced form (R). C) R from the generator has traveled over to the collectors and is reoxidized to O. D) A picture of a representative electrode array, 25 micron width electrodes, with a 25 micron separation, and 2 mm long. (Two of the array members on this array have been platinized.)





**Figure 2:** Chronoamperometric transients for the ETOF experiment in ferricyanide: the ferricyanide reduction current at the generator (red curve and axis) and the ferrocyanide reoxidation current at the collector (blue curve and axis). The generator electrode is briefly polarized, after this point the oxidative collector current increases until it reaches a maxima. The time between these two points is the time of maximum collection.

This type of experiment was first reported by Royce Murray, et al., to examine electron diffusion rates through conducting polymers that were inserted between two fingers on an electrode array(27). The two electrodes are within micrometers of each other and are often members of a microelectrode array (25, 28-33). The time it takes the product from the generator to diffuse to the nearest edge of the collector is the time of flight. If a potential pulse is applied to the generator, a burst of product diffuses to the collector (Figure 2), and the time of maximum collection,  $t_{mc}$ , is measured. The diffusion coefficient in of a Red/Ox species can then be calculated using Equation 3.

$$d = K\sqrt{Dt_{mc}}$$

Equation 3

Where  $d$  is the distance from the generator to the nearest edge of the collector,  $K$  is a geometric constant for the electrode system based on the height of the diffusion layer, the width of the electrodes and the gap between them(28),  $D$  is the diffusion coefficient. ETOF has the potential to greatly simplify the determination of diffusion coefficients as the concentration of the analyte, the area of the electrode, the viscosity of the solution, and the electron transfer kinetics can remain unknown. This equation is the general form as presented by Amatore (28) of an equation that was determined empirically by the Wrighton group for use in modeling diffusion of electro-generated species between electrodes in an array (33).

Another experiment similar to ETOF is the scanning electrochemical microscopy (SECM) experiment. This uses a ultramicroelectrode (UME) electrode held a short distance vertically from a surface. The surface, or substrate, and the UME, or tip, are used as a generator and a collector electrode and typically biased so that opposing reactions are happening at each electrode(34-39). SECM is typically used for surface imaging of electrode arrays and thin film studies, which both make use of the fact that the tip can be moved and gives spatial data about the resistive or conductive nature of the substrate. It can also be used for kinetics studies and the determination of diffusion coefficients(34, 36, 39). This is done by determining the diffusion coefficient ratio of the couple in the solution from the limiting currents of the collector and from the feedback to the generator. By knowing the current at time infinity, the electroactive area of the electrode, and the starting concentration of one of the analytes, one can determine the diffusion coefficients for both members of the couple(35). This experiment returns to a methodology that requires knowledge of the electroactive area of the electrode, and the initial concentration of one of the species, however longer requires knowledge of how far the electrodes are separated.

The previous literature concerning ETOF is mostly theoretical work, modeling the diffusion of a single molecule and then comparing the model to empirical data, where distance  $d$  was varied and the time of maximum collection measured at each distance. The geometric constant,  $K$ , was determined for a single molecule but never empirically proved to be a constant; especially for the case of multiple molecules diffusing in different buffer solutions. A key contribution in this paper is the rearrangement of Equation 3

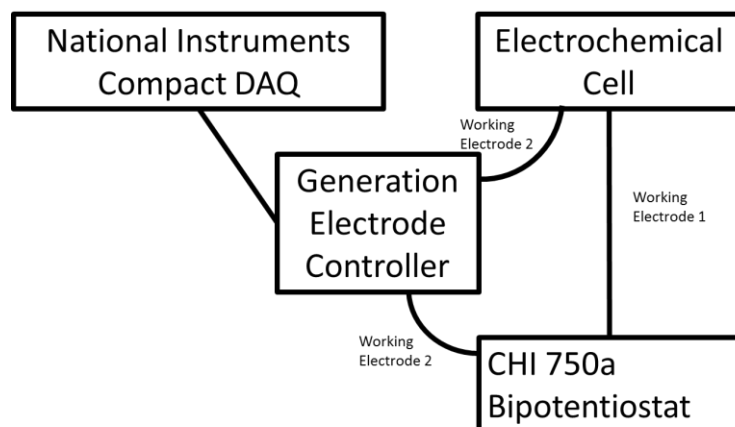
into Equation 4, as an alternative data analysis treatment for ETOF data.

$$\sqrt{t_{mc}} = \frac{d}{K\sqrt{D}} + B \quad \text{Equation 4}$$

If  $K$  is indeed constant, and  $d$  is held constant, then by selecting molecules with various and known diffusion coefficients, it would be possible to construct a curve with slope,  $\frac{d}{K}$ , and with an intercept  $B$ , which is a consequence of the x-intercept which occurs at the fastest diffusion rate able to be differentiated from  $t_{mc} = 0$  s. This “calibration curve” for the geometry of the array could then be used to determine the diffusion coefficient for any molecule in any solution. This experiment can be done without experiencing the loss of signal that occurs when  $d$  increases while performing a typical multiple distance experiment.

## Experimental

**Generator Controller and LabView Software:** Amatore used a multistat (Autolab Pgstat 20 and GPES

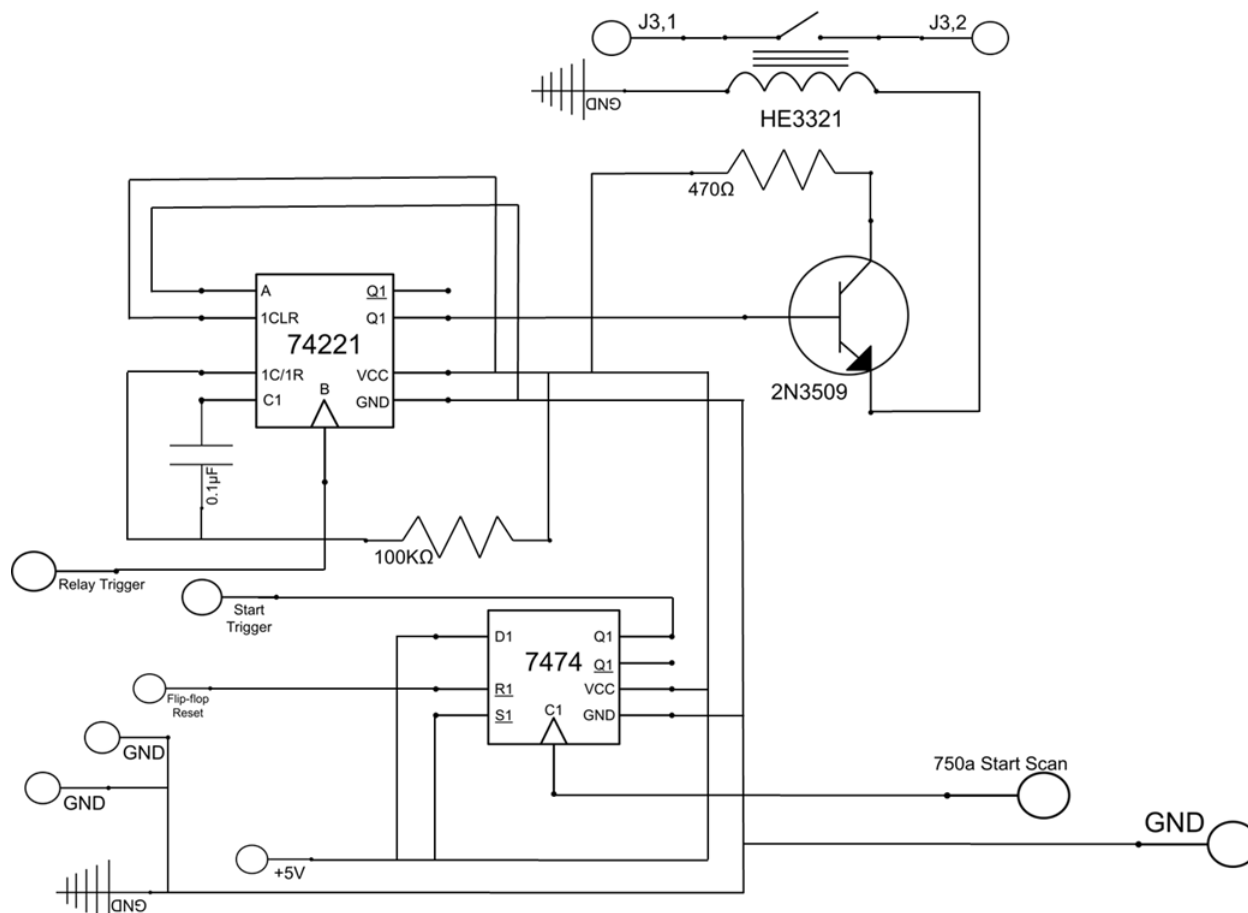


**Figure 3:** The connections between the Generation Electrode Controller (GEC), the National Instruments cDAQ modules, and the 750a potentiostat to produce a pulsed potential waveform at the second working electrode that would otherwise be impossible with the CHI 750a alone.

software from Ecochemie, Metrohm Switzerland) to perform these experiments. There are commercial instruments with the needed capability (Bio-Logic, Grenoble, France), but they are expensive. The bipotentiostats available in our lab, were not capable of leaving the second working (generator) electrode open circuit, nor of providing a potential pulse. Our solution was to

modify our existing bipotentiostat. The second working electrode from a CHI 750 potentiostat (CH Instruments, Austin, Texas, USA) was used to provide the potential to the generating electrode. As with most commercially available bipotentiostats the second working electrode only provides a static potential that is applied when the experiment initiates.

For the application described here, the generator is at open-circuit save for a brief potential pulse. To accomplish this, a relay was spliced into the second working electrode lead (Figure 3 Generating Electrode Controller, GEC). Timing and control of the relay was established by LabView software, and a



**Figure 4:** The circuit diagram of the generation electrode controller (GEC). The relay trigger, the start trigger, and the Flip-Flop Restart are connecting the controller to the NI 9403 Digital I/O module; the 750a Start Scan connects from the GEC to the cell control port on the CHI 750a; J3 is connected to the second working electrode running through the relay, and +5V and the ground at the lower right is from the power source for the system.

National Instruments CompactDAQ 9417 controller using a 9403 digital I/O module (National Instruments, Austin, Texas, USA), attached to a laptop computer (Figure 3). A circuit board was constructed to connect the relay and control and capture the digital signals between the potentiostat and the digital I/O module. The GEC went through several iterations before arriving at a final design shown in Figure 4. The D-flip-flop captures the downward “start-scan” pulse from the CHI 750a, latches the prompt so it will not be missed by the LabView software looping until the digital I/O line associated with the start trigger changes

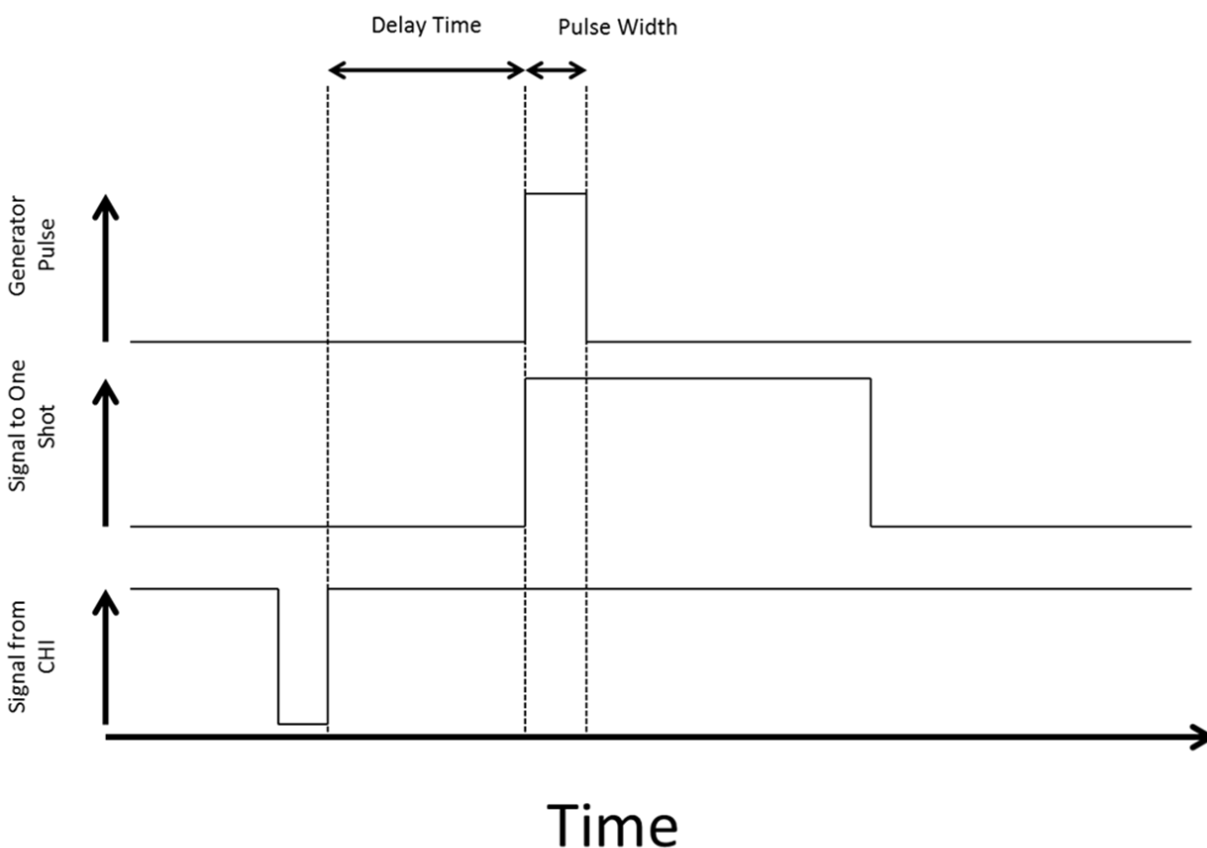
state. A software timer (seconds) selected by operator input, starts and at time-out, a second I/O output line, relay trigger, triggers the one-shot. The one-shot energizes the relay for a brief 15 ms period, momentarily applying a potential to the generator. The NPN transistor at the Q output of the one-shot, shown in Figure 4, provides sufficient drive to close the relay. Setting the generator pulse by the one-shot's RC time constant ensures that generator pulses are always a constant free from software timer and I/O uncertainties. The combination of the GEC and LabView software allows the collection of 10 sets of 10 generate-detect replicates, signal averaged over about 25 minutes.

**Working Electrode Array:** The electrode array was fabricated at the University's HiDec microfabrication facility. An array of 16 individually addressable microelectrodes, 25  $\mu\text{m}$  by 2 mm gold band working electrodes with a 25  $\mu\text{m}$  separation were constructed using standard photolithography procedures on silicon wafers, then diced. In ETOF experiments, one micro-band electrode served as the generator, with flanking bands (2), serving as the collector electrodes.

**Chemicals:** 5 mM potassium ferrocyanide (Certified ACS grade, Fischer), 5 mM potassium ferricyanide (Certified ACS grade, Fischer), 5 mM ruthenium (III) hexamine chloride (highest purity available, Alfa Aesar), and 5 mM dopamine HCl (99%, Alfa Aesar) were obtained and used as received. All except the dopamine were prepared in 0.1 M KCl (ACS grade, J.T. Baker) as their supporting electrolyte. Dopamine (5 mM) was prepared in 0.1 M phosphate buffer at pH 7.2, using mono and dibasic forms of sodium phosphate (Aldrich).

**Experimental Parameters:** The potentials applied to the generating and collecting electrodes were determined from cyclic voltammograms (CV) of the Red/Ox analytes. CVs were taken over potential ranges for the analytes in question, using a three electrode system; the working electrodes were members of the array, a SCE as the reference, and platinum flag as the counter electrode. The potential window of  $-0.4\text{ V}$  to  $0.4\text{ V}$  vs SCE was used for ferricyanide, ferrocyanide and ruthenium (III) hexamine;  $-0.1\text{ V}$  to  $0.55\text{ V}$  for dopamine. Potentials applied to the generator and the collectors were such that diffusion limited anodic and cathodic currents were achieved at these electrodes. The ETOF parameters were as follows: for ferrocyanide: the generator was pulsed to  $-0.2\text{ V}$ , while the collectors were held at  $0.4\text{ V}$  in a solution of 5 mM ferricyanide; for ferricyanide: the generator pulsed to  $0.4\text{ V}$  and the collectors

held at  $-0.1$  V vs SCE in 5 mM ferrocyanide in 0.1 M KCl; for ruthenium (II) hexamine: the generator was pulsed to  $-0.4$  V, while holding the collectors at 0.1 V in 5mM ruthenium (III) hexamine; and for



**Figure 5:** Timing diagram (not to scale) of the circuit that applies the potential to the generating electrode, triggered by a downward pulse from the “start-scan” output of the CHI 750a. A software timer in the CompactDAQ 9417 provides a 2.5 s delay time and then sends a pulse from a NI9403 digital I/O module to trigger a one-shot that applies a 15 ms pulse, to the relay that applies the second working electrode potential.

dopamine: the generator was pulsed to 0.55 V while the collectors held at  $-0.1$  V.

The ETOF and our alternative ETOF experiments were performed amperometrically with the collector polarized for the duration of the experiment, 5 s, and a generator pulse of 15 ms within that period, Figure 5. Experiments were performed in a faraday cage to reduce environmental noise. Dopamine is oxygen-sensitive, so these experiments were performed in oxygen-purged solutions, held under nitrogen. One hundred generate-collect experiments were performed in sequence and the collector currents were signal averaged to remove white noise. The averaged data was used to determine  $t_{mc}$  for the diffusing species.

In a second set of experiments, the same analytes were used, but the electrolyte solutions were dissolved in 20% v/v ethylene glycol solution to increase the viscosity by approximately twofold.

## Results and Discussion

A glance at the Levich equation gives some insight into the burden of determining diffusion coefficients of electrochemically active species. To extract a diffusion coefficient from the slope of the Levich line, values for kinematic viscosity of the solution, area of the electrode, and the number of electrons transferred must be known. Other methods, based on migration, also require known values for properties of the solution which must either exist in the literature or be measured separately. If the Red/Ox species is placed in a different buffer with a different viscosity,  $D$  will change along with all the other solution parameters. The advantage of the micro-electrode array method is that Equation 4 only depends on the geometry of the array and is independent of solution parameters.

**Applied potentials:** The electrode arrangement for the ETOF and our alternative experiments is shown in Figure 1. A chronoamperometric experiment is done at both the generator and collecting electrodes; the potentials can be set such that the products from the generator can be collected at the flanking electrodes. Most bipotentiostats apply potentials simultaneously at both working electrodes. If the potentials are applied simultaneously, eventually the current response at the collector will reach a steady-state-plateau (40). The steady state approach has a very high collection efficiency, but the difficulty is determining the time it takes to reach the current plateau; that is the “time of flight”. An additional complication is “feedback” or redox cycling (41), occurring when the diffusion layers of the generator and collector overlap, and products from the collector “feedback”, or “recycle” back to the generator. These two issues were resolved by applying a pulsed potential at the generating electrode, and pulse experiments then became the norm (28). Pulsing the potential at the generator produces a transient-peak seen in the amperometric display at the collector, Figure 2. Using this peak-shape, it is much easier to determine the time of flight, and the brief duration of the pulse does not generate enough material to allow feedback. Pulse widths for the generator were selected after consulting a paper by Amatore (28), so that the potential pulse applied to the generator would be less than the shortest time,  $t_{min}$  (Equation 5) for the molecule to diffuse between the two electrodes (otherwise redox cycling could occur).

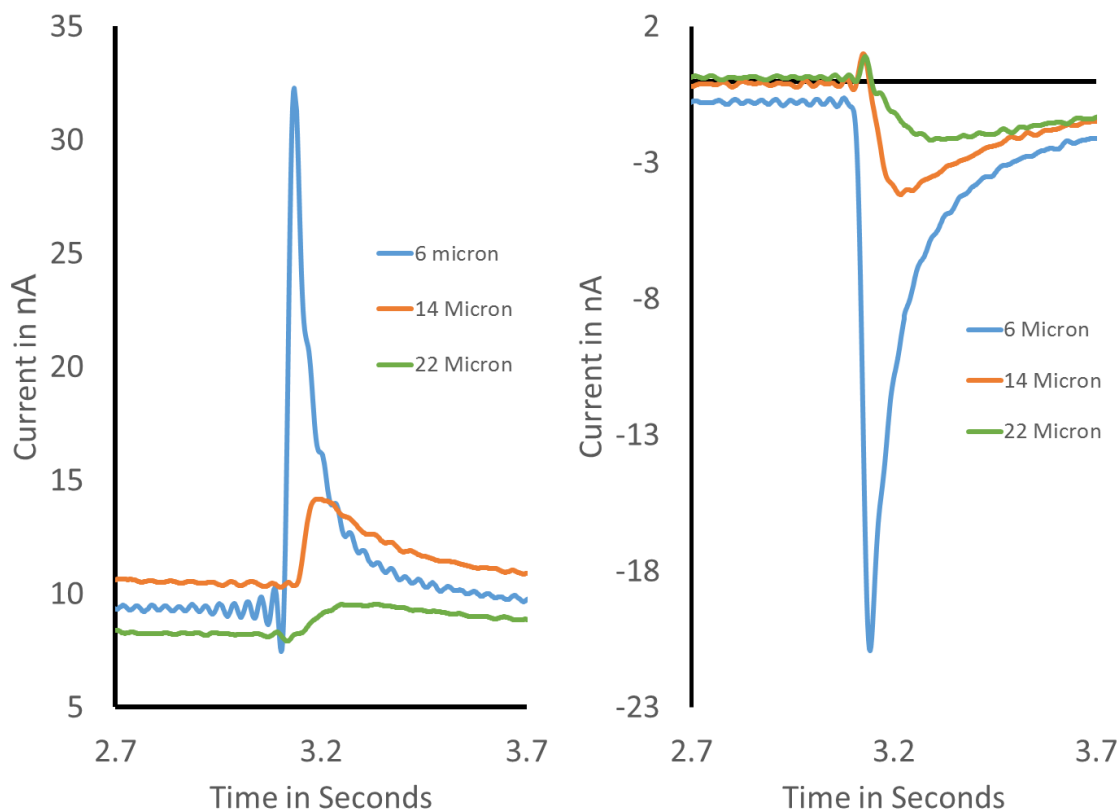
$$t_{min} = \frac{g^2}{4D}$$

Equation 5

In this equation  $g$  is the gap between the two electrodes and  $D$  is the diffusion coefficient. According to equation 5, the generator pulse was set to 15 ms.

**ETOF Data Analysis:** Pulse generation provides a peak in the collection current, allowing for easy determination of the  $t_{mc}$ . Figure 2 shows the transient current seen at the collector as a result of a potential pulse at the generator, and the time of maximum collection. The distance traveled has previously been measured in one of two ways: measuring from the edge of the generator electrode to the edge of the collector electrode, or from the center of the generator to the edge of the collector. We found that both methods gave comparable results, but used the edge to edge determination for  $d$ . The starting point,  $t_{mc} = 0$ , also has options: the rising edge of the generator pulse, the falling edge of the generator pulse, or the mid-point of the generator pulse. Any of these three options are valid as long as the defined start time is consistent throughout the experiment. For the case here, the time for maximum collection was measured from an artifact seen in the collector current, appearing as a spike the instant the generator is turned on. The current spike is caused by uncompensated resistance between the two electrodes (42, 43). The artifact eliminates the reliance on the temporal resolution of the second working electrode (2 ms), and allows time to be measured with the primary working electrode, which has a higher temporal resolution (1 ms). This was also useful, as it provides a point-in-time for synchronization, allowing the signal averaging of hundreds of repeat experiments, and provides a zero for the measurement of  $t_{mc}$ .



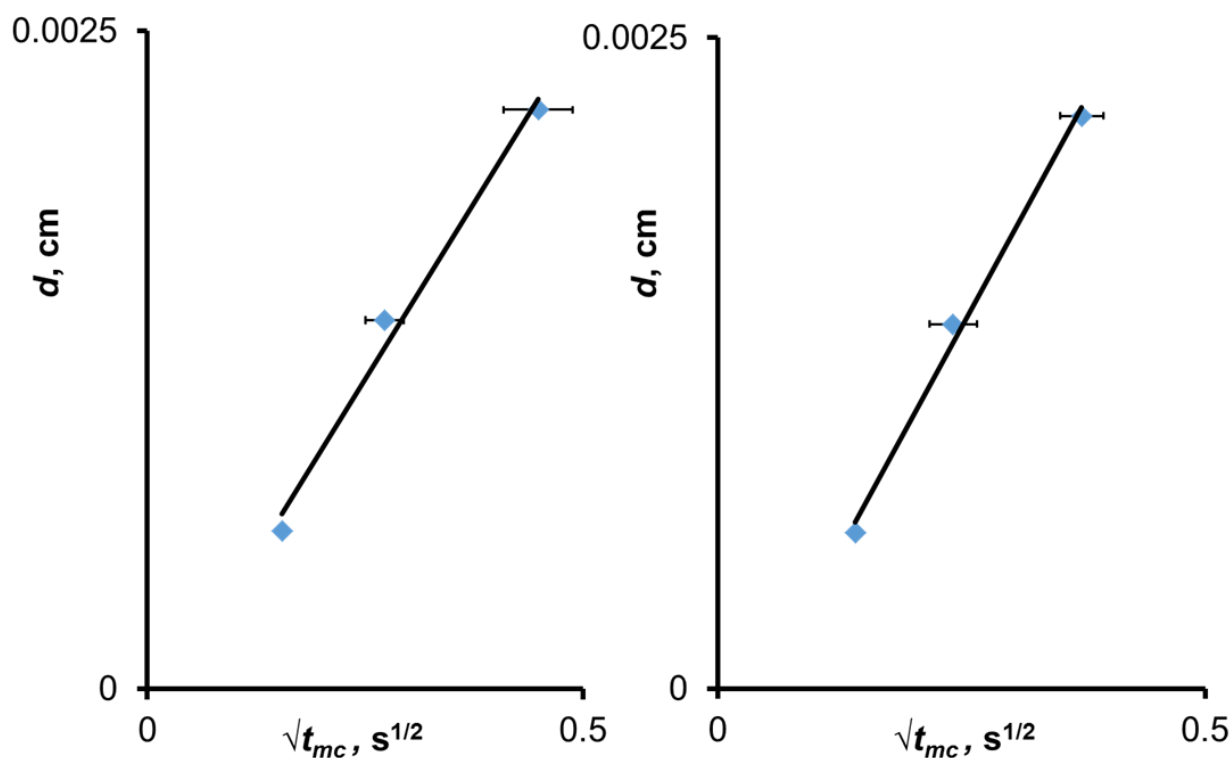


**Figure 6** : ETOF experiments in 5mM ferricyanide (left) and in 5mM ferrocyanide (right). On the left are the collector currents from the oxidation of generated ferrocyanide at three different distances from the generator (4 $\mu$ m electrodes with 4 $\mu$ m gap). On the right are the reductive collector currents on the same chip for generated ferricyanide. At time equals zero, the collector is polarized and double layer charging current is seen in the collector current. After the pulse at the generator at 3.11 s, a second signal is seen on the collectors as the generated species arrives. Both show data that has been smoothed using the CHI 750a's Fourier Transform smoothing option.

**Functionality of Hardware and Software:** To prove the functionality of the hardware and software, we chose an ETOF experiment from the literature where the diffusional distance,  $d$ , is varied by addressing three sets of flanking collectors in the array, each pair a larger distance from the generator. Band electrodes in the array were 4  $\mu\text{m}$  wide, 2 mm long, with 4  $\mu\text{m}$  separation. A solution of ferricyanide was used and the  $t_{\text{mc}}$  for the ferrocyanide generated measured at the three distances. Figure 6, left, shows an overlay of the collector currents at increasing distances from the generator. As the distance increases, more of the generated species is lost to the bulk solution resulting in decreased collection current. Because of the overall broadening of the peak, time of maximum collection is difficult to determine at larger distances, and the signal to noise decreases. Noise has not been recorded as a problem before, as larger band electrodes up to 2.0 cm long with micrometer separation provided larger currents (28).

The equation relating distance and time is given by Equation 3. This relationship, and the  $t_{\text{mc}}$  data from Figure 6, left, was used to construct Figure 7, left. From the slope of the line ( $K\sqrt{D}$ ), and a known diffusion coefficient for ferrocyanide (33) was used to determine  $K$ ,  $2.13 \pm 0.08$  ( $n=15$ ), which is a unitless parameter, for the array geometry used. This number was slightly larger (5%) than expected based on the theoretical work done by Amatore, indicating that the range for  $K$  should not exceed a value of 2, with array dimensions used here (28). We have not found ETOF experiments with similar geometries to ours, but we can say that our experimental  $K$  is very close to the theoretically predicted range.

In the previous experiment, the diffusive travel of ferrocyanide was used to determine  $K$ . With the geometric constant,  $K$ , known, it is then possible to determine the diffusion coefficients of Red/Ox species by simply measuring  $t_{\text{mc}}$ . As such, a second experiment was done to determine the diffusion coefficient for ferricyanide. A solution of ferrocyanide was used, and the  $t_{\text{mc}}$  of the ferricyanide generated was measured at the three distances shown in Figure 6, right. The diffusion coefficient for ferricyanide was determined to be  $7.3 \pm 0.7 \times 10^{-6} \text{ cm}^2/\text{s}$  ( $n=15$ ) using the slope of the line in Figure 7, right, and the previously determined  $K$  for the electrode geometry. This value is within 2% of the literature value of diffusion coefficient for ferricyanide (44). These results convinced us that our instrumentation was working correctly, and we proceeded with the development of using electrode arrays to determine diffusion coefficients.



**Figure 7:** Traditional ETOF treatment of the data shown in Figure 6. Left: Plot of the distance that the molecule traveled,  $d$ , as a function of the square root time of maximum collection,  $\sqrt{t_{mc}}$  for Ferrocyanide.  $K\sqrt{D}$  is the slope of the line so with a single analyte with known  $D$ ,  $K$  can be determined for a given electrode geometry.  $d = 0.0054(\pm 0.0002)\sqrt{t_{mc}} - 0.00029(\pm 0.00007)$ . Right: Plot of the distance vs the square root time of maximum collection for ferricyanide on the same electrode array. Knowing the  $K$  from the curve on the left one is then able to use the slope of this curve to determine the diffusion coefficient of ferricyanide.  $d = 0.0057(\pm 0.0001)\sqrt{t_{mc}} - 0.00014(\pm 0.00004)$ .

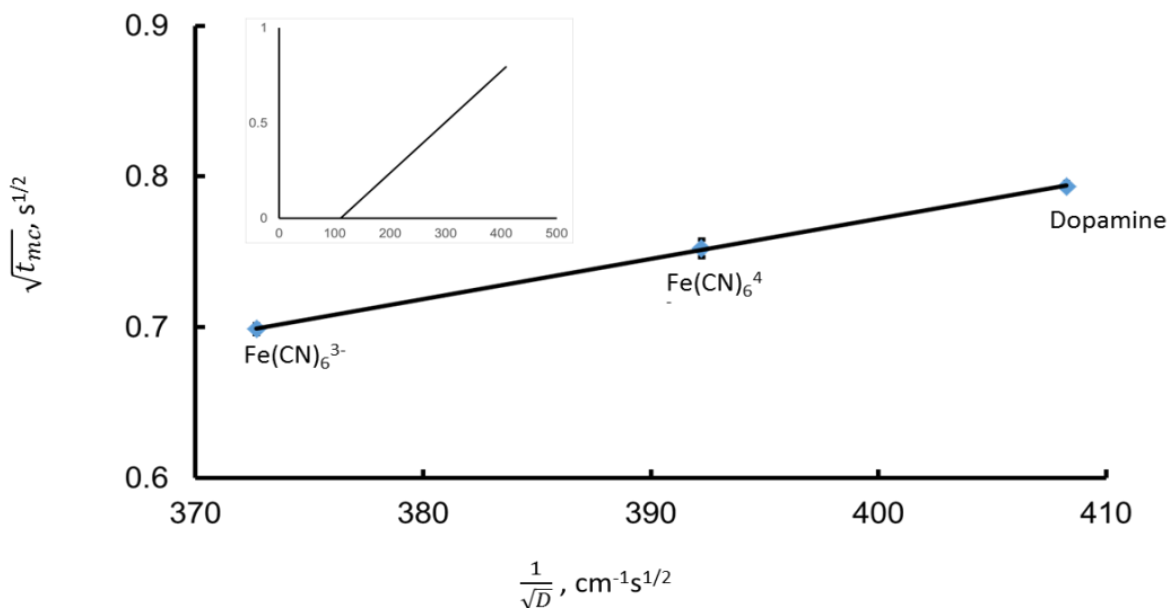
**Proving  $K$  a Constant for Multiple Red/OX Species:**  $K$  has been shown to be a constant when using a single redox species with a known  $D$ . The Wrighton group determined  $K$  for individual redox probe molecules in the early 1990s using the ETOF method with model compounds: in bulk solution with ruthenium (II) hexamine (33), in gels with ferrocene derivatives (29), and in solid polymers with silver ions that were stripped off of the generator electrode (45). Varga mathematically derived  $K$ , in a version of Equation 3 for a unique geometry describing glucose diffusion from a micropipette to an electrode, to determine glucose diffusion in solutions and gels (46). Slowinska heavily modified the ETOF technique to

study the capacitance of membranes on potentiometric sensors. The empirical data for hydrogen and silver ions was used to match computational models that were used to determine  $K$  (47, 48). Ky determined the diffusion of 4-hydroxy-(2,2,6,6-tetramethylpiperidine) in collagen matrices by first determining the geometric constant for the electrode geometry using the same probe molecule through solutions of various glucose concentrations, effectively increasing the viscosity; the work was still limited to a single analyte (31). Liu devised a method for the determination of diffusion coefficients using ETOF with cyclic voltammetry instead of the more traditional chronoamperometric methods; this was done by first modelling the cyclic voltammograms of ruthenium (II) hexamine and then verifying the resultant  $K$  with empirical data (25).

However, by only focusing on a single Red/Ox analyte, they did not challenge whether the  $K$  that they determined for their electrode geometry was applicable to other analytes using the same geometry. While other groups have focused on theory, no experimental data has been collected proving whether  $K$  is constant for multiple analytes.

The elegance of Equation 4 is in what is not found in the equation. There is no reliance of solution viscosity, number of electrons transferred, electron transfer kinetics, electrode rotation rate, nor solution conductivity. What follows is evidence that diffusion coefficients for multiple Red/Ox species can be determined quickly by using Equation 4, and that  $K$  is a constant for multiple Red/Ox species, in solutions of various viscosity. The line in Figure 8 was constructed by measuring  $t_{mc}$  for Red/Ox species with known diffusion coefficients (33, 44, 49) in a particular buffer system/solution, according to Equation 4. The analytes chosen were ferricyanide, ferrocyanide, and dopamine o-quinone (assumed to have the same diffusion coefficient as dopamine) to make the “calibration curve” for the array. The curve in Figure 8 can be approximated by a linear fit ( $\sqrt{t_{mc}} = 0.0027 \frac{1}{\sqrt{D}} - 0.2949$ ), with a slope that is equal to  $d/K$ . The distance between the electrodes is constant ( $d$ ) suggesting that the  $K$  is solely based on geometric parameters instead of being influenced by properties of the electrolyte solution, and it appears to be constant for multiple Red/Ox analytes. The y-intercept is meaningless but is a consequence of the fact that there must be an x-intercept. This is important because it means that at a given separation between the two electrodes,  $d$ , as the rate at which a species diffuses increases;  $t_{mc}$  must necessarily decrease.

Eventually  $t_{mc}$  will reach such a small value such that  $t_{mc}$  cannot be resolved from  $t_{mc} = 0$  s, given the temporal resolution of the data collection system. For the electrode separation used here, the x-intercept predicts that species diffusing faster than approximately  $1 \times 10^{-4}$  cm<sup>2</sup>/s cannot be determined. Most species in solution diffuse slower than this, so the technique is generally applicable for redox species in aqueous solution.



**Figure 8:** This curve plots the square root time of maximum collection,  $t_{mc}$ , as a function of known diffusion coefficients,  $D$ , reciprocal square root for various redox couples. The line can be used to determine the diffusion coefficients of multiple unknown species, just by measuring their time of maximum collection at a single collector distance. The fit of the line is  $\sqrt{t_{mc}} = 0.00267(\pm 0.00004) - 0.29(\pm 0.01)$ . The insert shows the plot extrapolated back to the origin, note the intercept of the abscissa represents the fastest diffusion coefficient that can be determined at an electrode

The line constructed from Equation 4 can also be used to determine the diffusion coefficients for Red/Ox analytes by simply measuring  $t_{mc}$ . The  $t_{mc}$  for ruthenium (II) hexamine chloride was measured as, 0.405 s, and using the calibration curve from Figure 8, the diffusion coefficient for ruthenium (II) hexamine was determined to be  $8.4 \pm 0.5 \times 10^{-6} \text{ cm}^2/\text{s}$ . This is because ruthenium (II) hexamine was generated from the ruthenium (III) hexamine in solution, and this method determines the diffusion coefficient of the diffusing species between the two electrodes. This determined diffusion coefficient was found to lie within 7% of the literature value and the null hypothesis could not be rejected at the 95% confidence interval, so it can be said that the value determined was equal to that found in the literature (33).

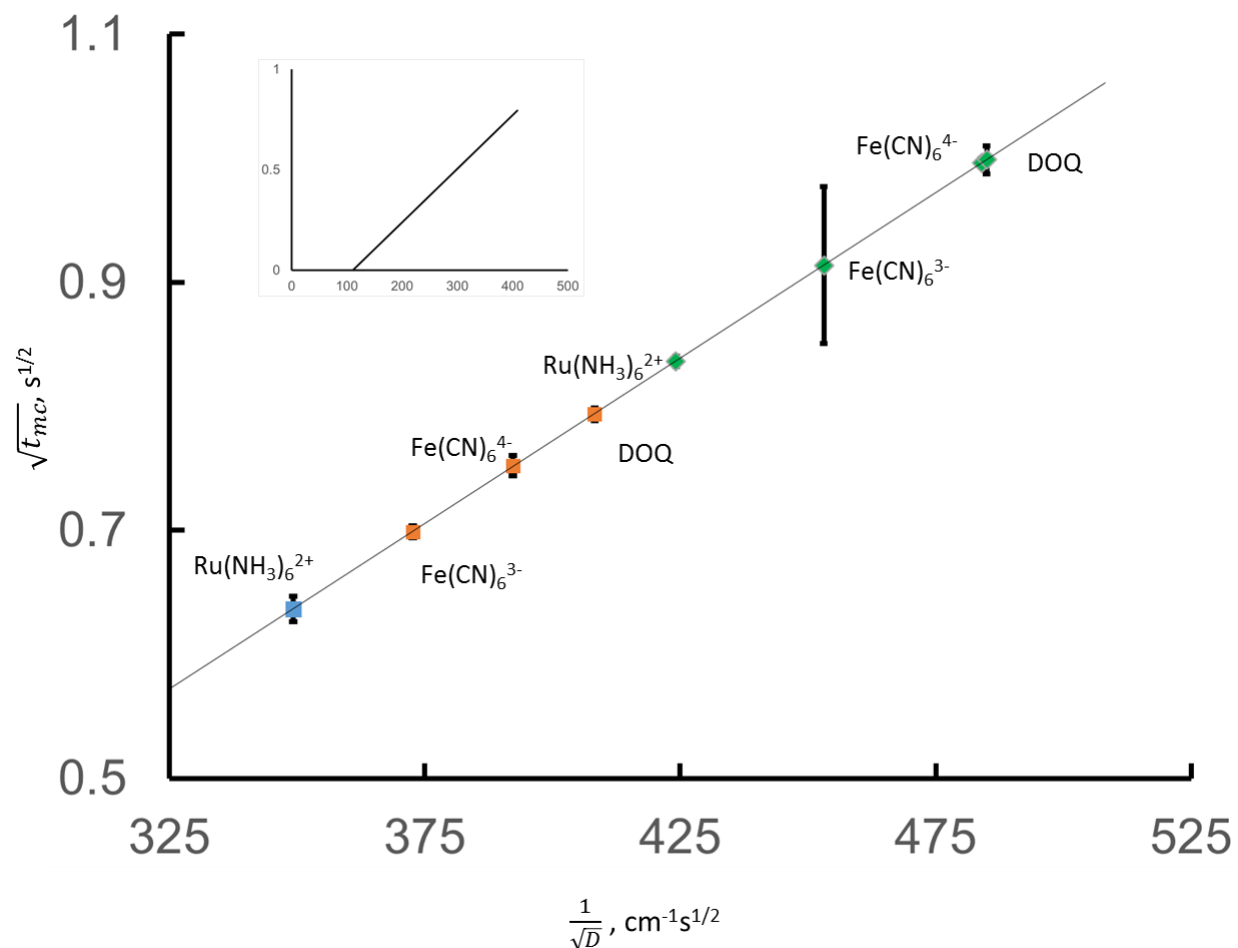
**Determination of Diffusion Coefficients in Viscous Solutions:** These results show that this kind of data treatment/calibration curve can be used for the rapid determination of the diffusion coefficients of

Table 1: Estimation of diffusion coefficients in solutions of increased viscosity using micro electrode arrays				
Species, generated	D, Literature, $\text{cm}^2/\text{sec} \times 10^6$	D, $\text{cm}^2/\text{sec} \times 10^6$ 20% v/v ethylene glycol/water	D estimate, Stokes-Einstein, $\text{cm}^2/\text{sec} \times 10^6$	% Difference
ferricyanide	7.20	$5.1 \pm 0.4$	4.45	13
ferrocyanide	6.40	$4.5 \pm 0.1$	4.01	11
ruthenium(II) hexamine	7.86	$5.8 \pm 0.2$	4.82	17
dopamine-orthoquinone	6.00	$4.4 \pm 0.1$	3.70	17

multiple Red/Ox species in solutions of varying composition. For example our alternative data analysis of the ETOF experiment for dopamine was performed in 0.1 M phosphate buffer, while the others used 0.1 M KCl as their supporting electrolyte.

We examined viscosity effects by performing additional experiments in solutions of 20% v/v ethylene glycol/water as the solvent for the supporting electrolyte. The time of maximum collection for ruthenium (II) hexamine, ferricyanide, ferrocyanide, and dopamine shifted to 0.70 s, 0.835 s, 0.993 s, and 0.998 s

respectively. The calibration curve in Figure 8 was used in Figure 9 to determine diffusion coefficients of each in the more viscous solutions (the green points in Figure 8). The diffusion coefficients determined for ferri/ferro cyanide, dopamine and ruthenium hexamine in the more viscous solution were determined from Figure 8 and are shown in Table 1.



**Figure 9:** The determination of the diffusion coefficients of ferricyanide, ferrocyanide, ruthenium(II) hexamine, and dopamine o-quinone (DOQ) through solutions with higher viscosity (20% v/v Ethylene glycol/water). The original aqueous calibration curve (orange markers), a determined diffusion coefficient for aqueous ruthenium (II) hexamine (blue marker), and the diffusion coefficients in increased viscosity (green markers), based on the calibration shown in Figure 8.

As expected the diffusion coefficient for each of the Red/Ox analytes has decreased by approximately 30% in the more viscous solution. Because there are no literature values for the diffusion coefficients in the more viscous solutions, we made estimates using the Stokes-Einstein equation: (50),

$$D = \frac{Tk_b}{6\pi r\eta}$$

Equation 6

Where  $D$  is the diffusion coefficient,  $T$  is the temperature,  $k_b$  is Boltzmann's Constant,  $r$  is the radius of the spherical particle, and  $\eta$  is the viscosity of the solution. Predicted values for the diffusion coefficients in the more viscous solutions are shown in the table, assuming that the change in solvent only affects the viscosity. Other than the Stokes-Einstein equation there has been little theoretical work predicting how diffusion coefficients should change with changes in viscosity. It has been shown that relative errors between experimental and the theoretical prediction from Stokes-Einstein range between 12 and 35% with more extreme errors possible based on the non-spherical nature of the particles that diffuse (50). Dopamine o-quinone and ferrocyanide in figure 8, appear to have the same diffusion coefficients in the 20% v/v ethylene glycol/water solution. Ferrocyanide has the expected slower diffusion coefficient in the more viscous solution, dopamine o-quinone however has a faster than expected diffusion coefficient. This could be because of unexpected solvent effects between the dopamine and the 20% v/v ethylene glycol solution altering the hydrated radius to smaller value. The determined diffusion coefficient for dopamine o-quinone is still in the low end of the expected range of disagreement (12-35%) with the Stokes-Einstein estimation.

## Conclusions

This alternative treatment of ETOF can quickly determine many diffusion coefficients for Red/Ox species from a single experiment. The method outlined here can be used to determine the diffusion coefficients, as viscosity increases, empirically and swiftly. This method can also be used to determine if diffusion coefficients predicted using the Stokes-Einstein equation make empirical sense. The percent error between the diffusion coefficients estimated with Stokes-Einstein and the diffusion coefficients determined empirically using our alternative method of ETOF was in the range of 11-17% error of the predicted values, and the acceptable range of percent errors is from 12-35%(44). These values fall within in the acceptable range, and were determined without expensive modeling software or time consuming experiments.

Future applications could include determining diffusion coefficients through membranes. Since diffusion



coefficients through membranes set the sensitivity of amperometric sensors, we envision, this research could be used to recalibrate implanted electrochemical sensors in biological applications, as well as a variety of other applications. There are as yet other possible applications that have not been considered because ETOF lacks the visibility needed for people to begin working on these important applications.

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**Optimization of Electrochemical Time of Flight Measurements for Precise Measurements of Diffusion Coefficients over a Wide Range in Various Media**

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## **Abstract**

Diffusion coefficients are an important physical property to both electrochemistry and the role they play in sensors, batteries, and catalysis. Electrochemical Time of Flight (ETOF) is a powerful and elegant technique that has largely been relegated to verifying diffusion modeling and theoretical work; instead of seeing prominence for its ease of use and broad applicability. ETOF does not require setting delicate hydrodynamic conditions that rotating ring disk (RDE) methods do, nor any mechanistic foreknowledge that is required about the electrochemical system, making this ideal for determining the diffusion coefficients for non-elucidated systems. Using experimental parameters optimized by computer modeling, an empirical calibration was generated for a platinum micro electrode array, and diffusion coefficients were determined for ferrocene, ruthenium(III) bisbipyridine dichloride, vanadyl acetoacetate, and tetrabutylammonium dioxovanadium(V) dipicolinate some of which were previously unreported.

## Introduction

In addition to being a fundamental property of any solute, diffusion coefficient values are crucial, especially in several applications of electrochemistry. Diffusion plays a pivotal role in mechanistic investigations as well as for amperometric and potentiometric sensors(1), electrosynthesis(2), redox flow batteries(3), etc. These are all applications where the electrical response and/or the global outcome depends on the diffusion of molecules to and from an electrode. Typically diffusion coefficients values,  $D$ , are determined either by electrochemical means using scanning electrochemical microscopy (SECM)(4-9), or by rotating ring disk (RDE) methods(10-17). RDE methods are based on the Levich equation, equation 1. The RDE limiting diffusion current plateau value,  $i_{Lev}$ , is related to  $D$  but also to the number of electrons transferred,  $n$ , the Faraday's constant,  $F$ , the active area of the electrode,  $A$ , the diffusion coefficient of the analyte (electroactive),  $D$ , the kinematic viscosity of the supporting electrolyte,  $\nu$ , the rotation rate,  $\omega$ , the analyte concentration of the in bulk,  $C_s$  (12). When  $i_{Lev}$  values are plotted as a function of  $\sqrt{\omega}$ , the slope containing multiple terms including the diffusion coefficient.

$$i_{Lev} = 0.620nFAD^{\frac{2}{3}}\omega^{\frac{1}{2}}\nu^{-\frac{1}{6}}C_s \quad \text{Equation 1}$$

The issue is that the other terms in the slope must be known to calculate the diffusion coefficient,  $n$ ,  $A$ ,  $\nu$ , and  $C_s$ . Acquiring values for these terms can be a little troublesome: such as the active surface area of the electrode, which could vary significantly from the geometric surface area; number of electrons transferred that, due to specific mechanistic circumstances, may be constant but not be an exact integer(18) or vary with the probed time range(19); kinematic viscosity of the solution; the concentration of the analyte. All of these can be determined through additional experiments performed with the required accuracy, but this makes for an additional burden. In addition, every electroactive analyte experiences different diffusion coefficient values in different media. Using RDE to determine diffusion coefficients requires a new Levich plot be constructed for each Red/Ox analyte every time any of the terms  $n$ ,  $A$ ,  $\nu$ , and  $C_s$  change, adding cumbersome details to the determination. Finally, it should be noted that equation (1) relates to the diffusion coefficient of the bulk electroactive species and not to its product.

A variant of this principle was also introduced to determine simultaneous values of  $n$  and  $D$  for any electrochemical wave through the combination of the limiting steady state current measured at an

ultramicroelectrode to the Cottrell current measured in the same solution at a millimetric electrode(19). However the method requires a high experimental precision and accuracy and should then be reserved to situations in which time-dependent kinetics may affect the then time-dependent  $n$  value.

In SECM, the probe serves as the detector electrode and is moved vertically to different distances from a surface that is generating a redox active species. The diffusion coefficient is determined by measuring the time of maximum collection at different distances from the surface. However such SECM-based method may not be applicable in structured electrolyte media such as polymers or sol-gel materials in which diffusion coefficient values may be required(17).

In fact, the very principle of such SECM experiments is similar to that of electrochemical time of flight (ETOF) ones as both involve generation and collection of the species of interest. Electrochemical time of flight (ETOF) was first developed by Royce Murray(20) to measure the time that a generated species took to travel from one electrode to an adjacent detecting electrode located at a known distance away; the travel time could then be related to the diffusion coefficient. Murray was interested in determining the apparent diffusion coefficients of electron-hopping across conducting polymers sandwiched between two electrodes (20). Hence, the SECM and ETOF approaches only differ from the fact that the former relies on collector-generator current feedback magnitude (viz., on collection efficiencies(21)) while the second one relies on measurements of travel times of the generated species of interest between its source (the generator electrode) and its sink (the collector electrode(s)). This travel time,  $t_{mc}$  is indeed related to the diffusion coefficient,  $D$ , through Einstein's equation, equation 2, the distance traveled or gap between the electrodes,  $g$ , and a numerical dimensionless constant,  $K$ , that depends exclusively on the electrode array geometry and not at all on the electrochemical medium.

$$g = K\sqrt{Dt_{mc}} \quad \text{Equation 2}$$

$K$  can be predicted by computational means for generator-collector systems having ideal geometries (22, 23) or evaluated empirically based on calibrations performed for micro-systems of any geometry(24-26). Most of the literature has been reporting ETOF applications as a way to verify computational models of diffusion between electrodes, viz., as a way to assess experimentally the "ideality" of the microfabricated

systems, but rarely as an extremely adequate and simple technique for determining diffusion coefficients. Previously, transit times of analytes at different distances between the generating and detecting electrodes were measured to determine diffusion coefficients with ETOF. These experiments have been limited to model compounds such as ruthenium hexamine(23, 25), ferricyanide(22, 23), and hydroxymethylferrocenium(27) for example, or the analyte of interest in a given study(24). The advantage of ETOF over other methods is that it does not require the foreknowledge all the solution parameters mentioned above, as required by RDE methods, nor does it require specialized equipment as SECM does. Following the advent of mass produced microelectrode arrays and despite the simplicity of the technique it is surprising that ETOF is not a standard method for determining unknown diffusion coefficients.

In our previous work, we presented a novel data treatment for electrochemical time of flight (ETOF) that we believe provided a simple and elegant method for determining diffusion coefficients. In this approach a calibration curve is constructed by measuring  $t_{mc}$  with a given micro-device (viz., for a single and constant distance,  $g$ , between the generator and collector electrodes) for electrogenerated species with known diffusion coefficients in the medium of interest. Rewriting equation 2 as follows.

$$\sqrt{t_{mc}} = \frac{g}{K\sqrt{D}} \quad \text{Equation 3}$$

Plotting  $\sqrt{t_{mc}}$  variations as a function of  $\frac{1}{\sqrt{D}}$  known values for species with reported diffusion coefficients, produces a calibration line that can then be used to determine diffusion coefficients of unknowns species regardless of other characteristics of the medium and of the species. In this paper we have expanded electrolyte solutions to include non-aqueous solvents. In addition, computer modeling was used to optimize the pulse widths for the generator, so that diffusion coefficients could be determined over a large dynamic range with an adequate precision and accuracy.

### **Theoretical optimization of precision and accuracy in ETOF measurements**

Generator-collector micro-devices may be used under steady state regimes(21, 23). Referring to Figure 10, it is clear relying on steady-state operation is not suitable since any information related to the time of flight of the species B formed at the generator is lost, being somewhat convoluted with kinetic effects into

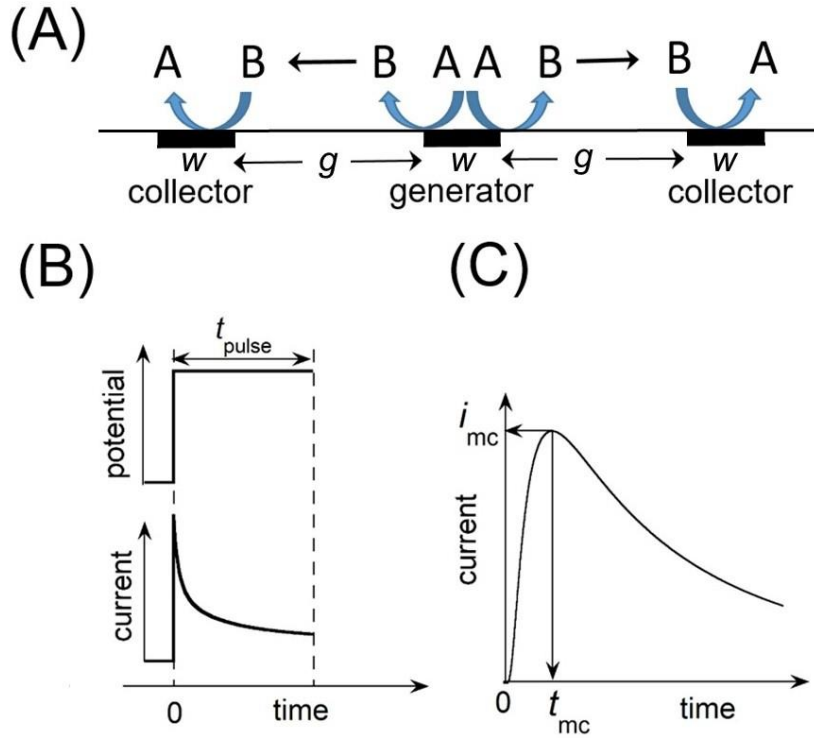


the common current intensities of the generator and collector electrode(s)(21). Conversely, applying a sufficiently brief potential pulse(18,24) to the generator electrode generates a local concentration of B in its very vicinity(28). Diffusional broadening of this concentration pulse over the insulating gap allows B to be captured by the collector(s). though, since the duration of the generator pulse is brief, B is rapidly exhausted. Hence, a peak-shaped current response is expected at the collector(22, 29) whose maximum time position,  $t_{mc}$ , is directly indicative of B shortest time-of-flight (TOF). Conversely, whenever the electrochemical detection of B regenerates A, this latter species may then diffuse back to the generator where it can regenerate B if the generator is still electrochemically active. This is bound to induce a redox-cycling mechanism while the generator may perform. As a consequence, the time position of the maximum collector current is deported towards large times since it now represents a convolution between the shortest TOF and the maximization of the redox-cycling efficiency. Therefore, in absence of any noise or background interference, the smallest is  $t_{pulse}$  value the most accurate is the application of equations 2 and 3. Conversely, the shortest is  $t_{pulse}$ , the smallest is the generated quantity of B and its electrochemical current at the collectors, hence the lowest is the precision of  $t_{mc}$  determination. This experimental conflict between accuracy and precision indicates that  $t_{pulse}$  value should be as close as possible from a maximum one,  $t_{pulse}^{max}$ , beyond which redox-cycling contributions will alter the ETOF measurements. However, it is understood that this value is not universal but depends on the rate of diffusion of B over the insulating gaps, viz., on  $t_{mc}$ .

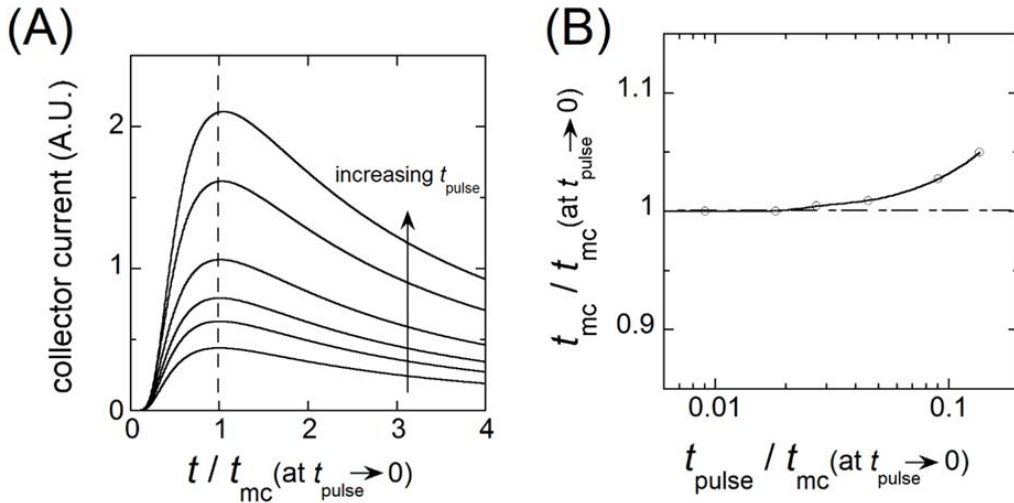
Therefore, though the present work is aimed to validate experimentally the principle of diffusion coefficients measurements by ETOF and determine experimentally the specific calibration curve relating  $D$  and  $t_{mc}$  values (equation 3) for any given micro-device irrespective of the electrolyte medium in which the measurements are performed, one needs first to estimate the relationship between  $t_{pulse}^{max}$  and  $t_{mc}$  values in order to guarantee both the precision and the accuracy of the ETOF calibration curve for the micro-device of interest. This is essential to ensure that the corresponding calibration curve may then safely be used for determining the unknown diffusion coefficients of other species in any media with the same micro-device. For this reason we wish to present the results of a simplified theoretical analysis aimed to evaluate the relationship between  $t_{pulse}^{max}$  and  $t_{mc}$  values.

This theoretical analysis was performed through Comsol simulations of ideal micro-device systems consisting of three parallel microband electrodes of common width,  $w$ , separated by insulating gaps of common width  $g$ , the whole assembly being imbedded flat in an insulating plane (Figure 10A)(21, 30). Owing to its scope, this theoretical analysis was simplified through considering that the electrodes present no vertical edges and that no metallic conductor was located within the gap separating the generator and collector electrodes in order to avoid theoretical complications related to any bipolar activity that may be induced over such isolated conductors(31, 32).

Initially, the semi-infinite bulk electrolyte is assumed to contain only a species A. This is reduced, or oxidized accordingly, at the generator electrode whose potential is stepped for a fixed time duration,  $t_{\text{pulse}}$ , at an adequate value and disconnected afterwards (Figure 10B). This generates a peak-shaped current response (Figure 10C) characterized by its current maximum,  $i_{\text{mc}}$ , and its time position,  $t_{\text{mc}}$ , after the beginning of the pulse ( $t = 0$ ).



**Figure 10:** (A) Schematic cross section of the device considered in the ETOF experiments.  $g$  is the gap distance between the electrodes selected to operate as generator and collector electrodes. These electrodes are indicated by black rectangles embedded in the insulating plane. (B) Schematic time dependence of the potential pulse applied at the generator and of the resulting generator current. (C) Corresponding collector response characterized by a peak-shaped current with its maximum characterized by its current, time coordinates  $(i_{mc}, t_{mc})$ .



**Figure 11:** Illustration of the influence of  $t_{pulse}$  on the position of the collector peak response,  $t_{mc}$  (A) Collector current response for increasing  $t_{pulse}$  values.  $t_{mc}$  has a constant value, independent of  $t_{pulse}$  when  $t_{pulse} \rightarrow 0$ . (B) Variation of the ratio  $t_{mc} / (t_{mc} \text{ at } t_{pulse} \rightarrow 0)$  as a function of  $t_{pulse} / (t_{mc} \text{ at } t_{pulse} \rightarrow 0)$ .

Figure 11A illustrates the outcome of the corresponding simulations by evidencing how the collector peak-shaped response vary when  $t_{\text{pulse}}$  increases. Figure 11B reports the variations of  $t_{\text{mc}}$  normalized to its ideal limit when  $t_{\text{pulse}} \rightarrow 0$ , as a function of  $t_{\text{pulse}}$  also normalized to the same limit. This dimensionless presentation of the theoretical results makes that Figure 11B is independent of the exact nature of the redox system A/B considered and of the geometry of micro-device used provided that its geometric characteristics approach closely the ideal ones defined above. Also, since the generator and collector potentials may be experimentally chosen so that the reduction of A, or its oxidation accordingly, and the detection of B are performed at the current plateaus of the corresponding electrochemical waves, Figure 11B is independent of the exact mechanisms involved during these electrochemical processes provided that species B is the one that diffuses over the micro-device gaps. In other words, would B be a follow-up product resulting from the chemical evolution of  $A^\pm$ , the primary intermediate formed by the reduction of A, or its oxidation accordingly, the half-life time,  $t^{1/2}$ , of the formation of B from  $A^\pm$  must be negligible vs.  $t_{\text{mc}}$ . Would that not be the case, the ETOF method would provide a “mixed” diffusion coefficient convoluting those of B and  $A^\pm$ . In order to guarantee that this is not the case, the ETOF method should be reserved to the analysis of redox systems that are not necessarily reversible but whose CV simultaneously displays the voltammetric waves of A (forward scan) and B (backward scan) and no intermediate wave within a single voltammetric scan performed at a scan rate,  $v$ , matching  $t_{\text{mc}}$ , viz. such as  $v_{\text{mc}} \sim RT/(Ft_{\text{mc}})$  where  $R$  is the perfect gas constant,  $F$  the Faraday number, and  $T$  the absolute temperature(33).

A first conclusion of this simplified theoretical analysis described in this section, shows that the optimization of its accuracy and precision imposes (Figure 11B) that the duration of the generator pulse duration,  $t_{\text{pulse}}$ , does not exceed a maximum value being ca. 0.03 to 0.04 times that of the experimentally determined  $t_{\text{mc}}$ . The validity of this criterion must therefore be verified after a given determination of  $t_{\text{mc}}$  in order to ensure the validity of the correspondingly determined diffusion coefficient value. A second conclusion is that the ETOF method applies to determine the diffusion coefficient of any species that is the primary intermediate of the reduction or oxidation of a precursor species, A, or is generated from this primary intermediate within a time-scale  $t^{1/2}$  that is much shorter than  $t_{\text{mc}}$ . The exact relationship between the maximum value,  $(t^{1/2})_{\text{max}}$ , of  $t^{1/2}$  and  $t_{\text{mc}}$  is evidently depending on the mechanism relating B to the primary intermediate  $A^\pm$ , but the use of equation 3 should provide correct results when  $t^{1/2} < t_{\text{mc}}$ .

## Experimental

**Hardware:** All electrochemical measurements were performed using a CHInstruments 750a bipotentiostat, and an in-house built circuit to control the generator electrodes. The in-house circuit was controlled by National Instruments LabVIEW software. The micro band electrode arrays of 16 platinum or gold fingers, 2 mm long, 25  $\mu\text{m}$  wide with gaps of 25  $\mu\text{m}$  were fabricated at the Arkansas High Density Electronics Center(HiDEC). Using either one of these arrays provided essentially identical results. An Ag/AgCl in saturated KCl reference electrode was used in aqueous solutions. For use in organic solvents a bridge, made with TBAPF<sub>6</sub> in acetonitrile was constructed over a SCE.

**Chemicals:** Potassium ferrocyanide (Certified ACS grade, Fischer), potassium ferricyanide (Certified ACS grade, Fischer), ruthenium (III) hexamine chloride (highest purity available, Alfa Aesar), ferrocene (98% Sigma Aldrich), ferrocenyl acetic acid (98% Sigma Aldrich), ruthenium bisbipyridine dichloride (Ru(bpy)<sub>2</sub>Cl<sub>2</sub>) (97% Sigma Aldrich) were all obtained and used as received unknowns were all obtained and used as received to make 5 mM solutions in the chosen electrolyte. Two solutions were considered. One (aqueous) consisted of 0.1 M KCl (ACS grade, J.T. Baker). The other (organic) consisted of 0.1 M TBAPF<sub>6</sub> (98%, Sigma Aldrich) in acetonitrile (HPLC Grade, Fischer).

The vanadyl acetylacetonate (VO(acac)<sub>2</sub>) was synthesized from Vanadium(V) Oxide (98%, Sigma Aldrich) and acetylacetone (Reagent Plus grade, Sigma Aldrich) according to the method below and was recrystallized for purification in dichloromethane(34). Note however that VO(acac)<sub>2</sub> oxidation has been reported to provide very rapidly VO(acac)<sub>3</sub>, especially if the solution contains already traces of VO(acac)<sub>3</sub>(35). The presence of VO(acac)<sub>3</sub> traces in our synthesized VO(acac)<sub>2</sub> samples was confirmed by comparing the infra-red spectra of the solution of our synthesized VO(acac)<sub>2</sub> with that reported for a pure VO(acac)<sub>2</sub> solution (Sigma Aldrich).

The tetrabutylammonium dioxovanadium(V) dipicolinate (VO<sub>2</sub>(dipic)) synthesis proceeded as follows. A tetrabutylammonium hydroxide ([n-Bu<sub>4</sub>N]OH) solution was prepared by diluting 100mL of 40% wt. in H<sub>2</sub>O (Aldrich) to 1 liter with deionized water to yield a 0.4M solution. Approximately 4.0g of solid V<sub>2</sub>O<sub>5</sub> ( $\geq 98\%$ , Aldrich) was added to 200mL 0.4M [n-Bu<sub>4</sub>N]OH in a 500mL Erlenmeyer flask and left to stir for approximately 18 hours at room temperature. The solution changed from a transparent yellow to nearly

colorless liquid with a very small amount of colorless solid floating on the top. The solid was vacuum-filtered from the solution using a Buchner funnel and the remaining filtrate transferred to a 250mL round bottom flask. The filtrate was then dried on a rotary evaporator for 1-2 hours at 46° C until the appearance of viscous, light-brown oil. The round bottom flask was placed on a Schlenk line and left overnight to evaporate to complete dryness, yielding a crude, white solid. The flask was then back-filled with nitrogen, sealed, and transferred to a glovebox under a nitrogen atmosphere for extraction of the air-sensitive intermediate, [n-Bu<sub>4</sub>N][VO<sub>3</sub>](36).

To complete the synthesis of VO<sub>2</sub>(dipic), approximately 1 gram of the [n-Bu<sub>4</sub>N][VO<sub>3</sub>], 0.5g 2,6-pyridine dicarboxylic acid (99%, Aldrich), and 3g 4Å molecular sieves were placed in 15-20mL dichloromethane (99.8%, anhydrous, Sigma-Aldrich) and left to stir for ~45 minutes. The sieves and insoluble solid were vacuum-filtered using a fritted funnel and flask. The remaining filtrate was then pumped down to ~4mL, layered with diethyl ether (dried, Sigma-Aldrich), and placed in the glovebox freezer for crystallization. After seven days the mother liquor was decanted and the wet crystals dried under vacuum. The crystalline solid was then transferred to a clean vial, triturated with fresh diethyl ether, dried again under vacuum and then weighed. Typical final product yield: 0.88-0.92 grams(37) of VO<sub>2</sub>(dipic).

#### **Conditions of electrochemical measurements:**

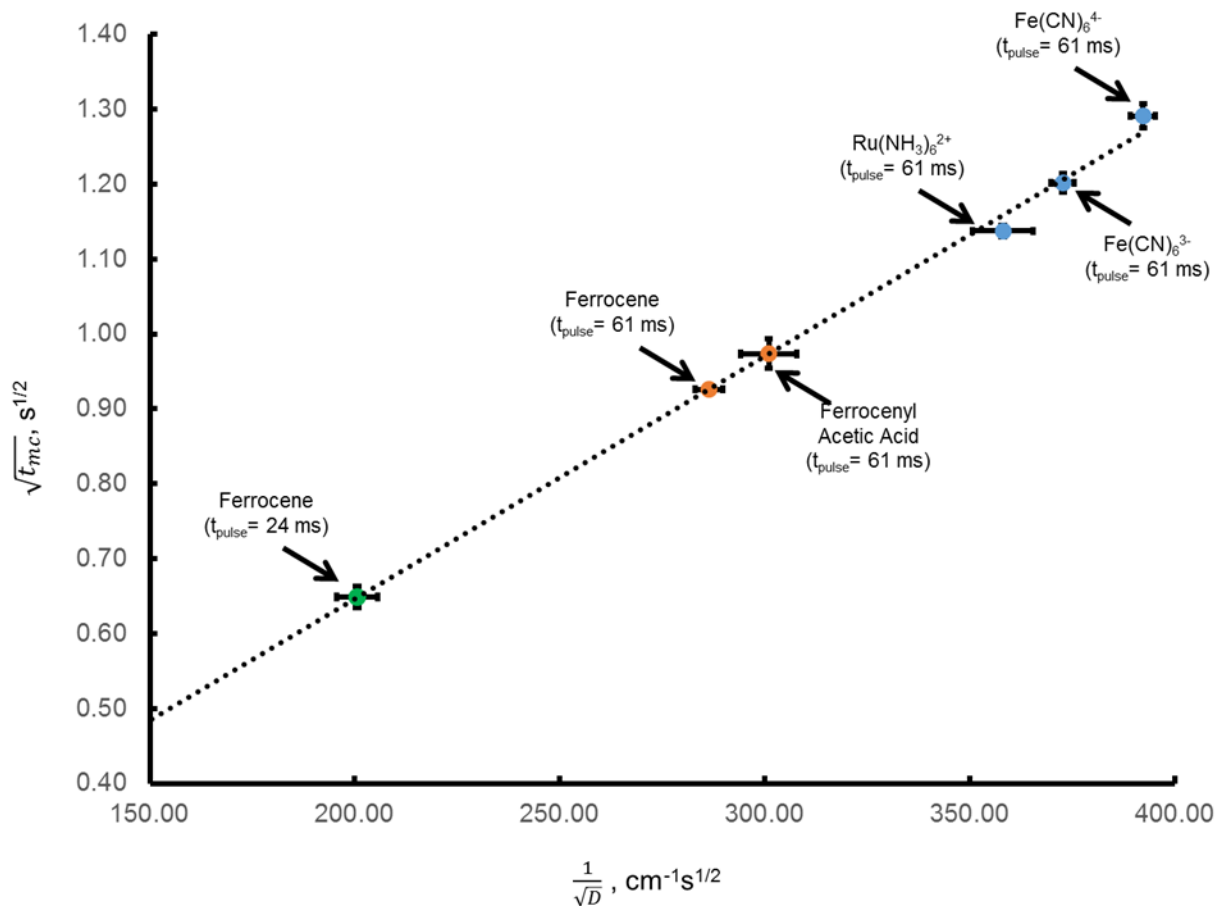
Three equally spaced (gap distance of  $g=75\mu\text{m}$ ) microband electrodes ( $w=25\mu\text{m}$ ) was selected by pairing sets of microbands available in the microfabricated arrays to perform ETOF measurements. The other 13 microbands present in the array were not connected. The generator electrode potential was imposed during a time duration,  $t_{\text{pulse}}$ , and the generator immediately disconnected to minimize current feedback. The two collector electrodes remained continuously poised(30).

The applied potentials values were determined from CVs of reported for the A/B couples performed in 0.1 M KCl or in 0.1 M TBAPF<sub>6</sub> / acetonitrile. For the latter case, unless otherwise noted, all potentials were measured versus a SCE with a double bridge. The potential window for the CV for ferrocene and ferrocenyl acetic acid solutions was from -0.6 V to +0.8 V, and the ETOF potentials were selected as follows: for ferrocenium, the generator was pulsed to +0.8 V while the collectors were held at +0.1 V. For

experiments with ferrocenyl acetic acid, the generator was pulsed to +0.5 V while the collectors were held at 0.1 V. The CV potential window for  $\text{Ru(II)(bpy)}_2\text{Cl}_2$  was +0.5 V to -0.2 V and the ETOF potentials set at +0.5 V for the generator pulse and -0.1 V for the collectors. The ferrocene/ferrocenium had a known diffusion coefficient in the literature(38) and  $\text{Ru(II)(bpy)}_2\text{Cl}_2$  was compared to an estimated diffusion based on its mass(39).

$\text{VO}_2(\text{dipic})$  electrochemical reduction forms a dimer when there are proton donating buffers such as tosic acid. So, the diffusion coefficient of the reduced form of this dimer was determined in acetonitrile with 0.1 M electrolyte  $\text{TBAPF}_6$  and the tosic acid buffer 0.0150M. The  $\text{VO}_2(\text{dipic}) / \text{VO}_2(\text{dipic})$  dimer couple is not a reversible redox one as the above ones. The voltammetric potential window was then extended over the range +1.6 V to -1.6 V vs a silver cryptand reference to identify the  $\text{VO}_2(\text{dipic})$  reduction wave and the oxidation one of the  $\text{VO}_2(\text{dipic})$  dimer. This enabled selecting ETOF potentials as +0.4 V for the generator pulse and +0.9 V for the collectors. The diffusion coefficient of the dimer was determined experimentally using ETOF and compared with those of the dimer determined by other by using Koutecky-Levich plots.

The mechanism undergoing at the  $\text{V}^{\text{IV}}\text{O}(\text{acac})_2$  oxidation wave is even more complex, leading to  $\text{V}^{\text{V}}\text{O}(\text{acac})_3$  (see below in Results and Discussion). The potential window for the CV of  $\text{VO}(\text{acac})_2$  was thus extended from +2.0 V to -2.0 V to identify the ETOF potentials. These were set at +1.0 V for the generator pulse and -1.5 V for the collectors in order to determine the diffusion coefficient of  $\text{VO}(\text{acac})_3$ .



**Figure 12:** Calibration curve constructed using standards compounds ferrocyanide, ferricyanide, and ruthenium (II) hexamine in 0.1 M KCl aqueous solution (blue points) based on their  $t_{mc}$  values determined from ETOF experiments using  $t_{\text{pulse}} = 61\text{ ms}$  and known diffusion coefficients. The dotted line represents the equation  $\sqrt{t_{mc}} = 3.23(\pm 0.03) \times 10^{-3} / \sqrt{D_B}$ . The incorrect  $t_{mc}$  values determined for ferrocene and ferrocenyl acetic acid in acetonitrile using  $t_{\text{pulse}} = 61\text{ ms}$  are shown by orange circles while the correct one measured for ferrocene using  $t_{\text{pulse}} = 24\text{ ms}$  is shown by a green point.

## Results and Discussion

An experimental ETOF calibration curve (Figure 12) was determined as reported before(30) using a constant value of  $t_{\text{pulse}} = 61\text{ ms}$  using either gold or platinum band electrodes. The calibration curve was constructed based on the experimental  $\sqrt{t_{mc}}$  values determined in 0.1 M KCl for ferrocyanide ( $6.40 \times 10^{-6} \text{ cm}^2/\text{s}$ )(23), ferricyanide ( $7.20 \times 10^{-6} \text{ cm}^2/\text{s}$ )(40) and ruthenium(II) hexamine ( $7.86 \times 10^{-6} \text{ cm}^2/\text{s}$ )(23) whose  $D$  values are reported in the literature. The regression line,  $\sqrt{t_{mc}} = 3.23(\pm 0.03) \times 10^{-3} / \sqrt{D_B}$ , was determined by enforcing it to pass through the origin in order to comply with the theory (18). From equation 3 and  $g=75\mu\text{m}$ , this slope provided  $K = 2.32(\pm 0.02)$ . It was also verified that for the three species



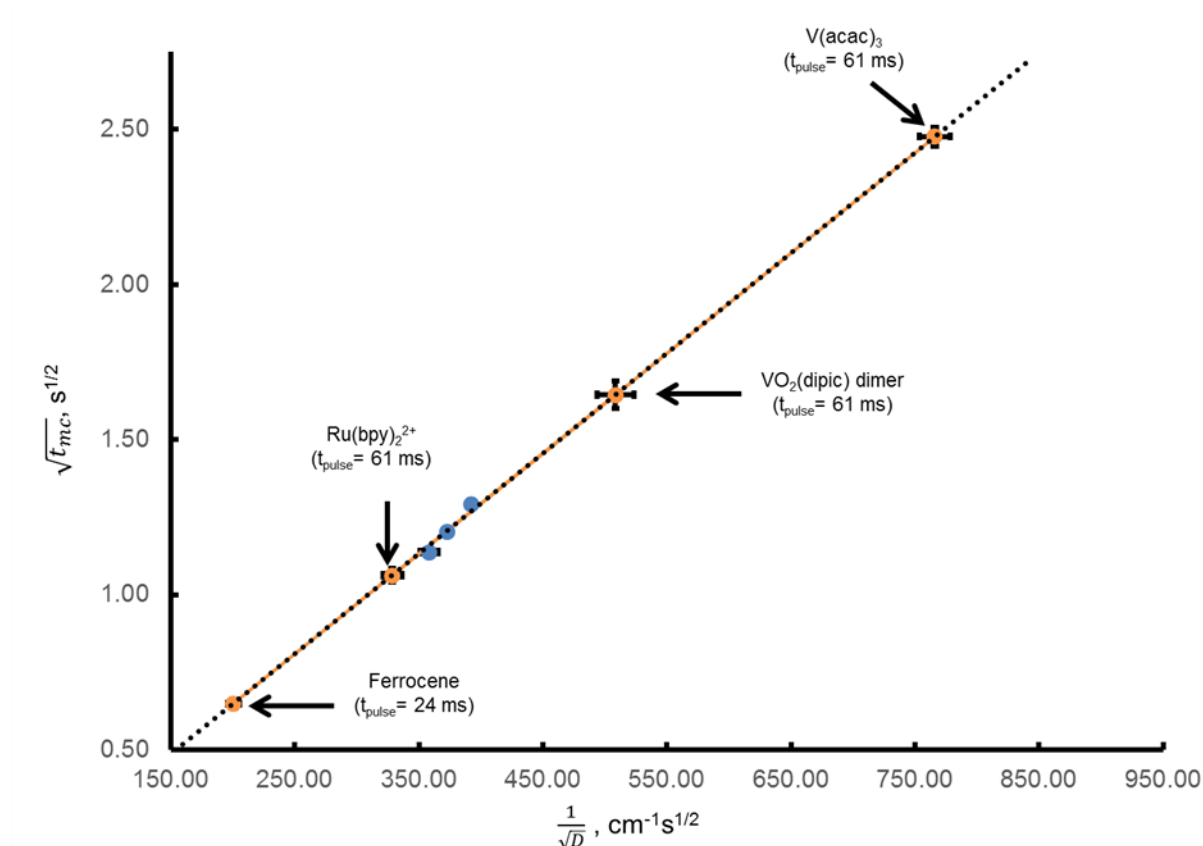
used as calibrants, using  $t_{\text{pulse}} = 61$  ms satisfied the criterion developed above. Indeed, based on the experimental  $t_{\text{mc}}$  values, one had  $t_{\text{pulse}}/t_{\text{mc}} = 0.037$  (ferrocyanide), 0.042 (ferricyanide) and 0.046 (ruthenium(II) hexamine), which is satisfactory owing to the theoretical results in Figure 11B.

Using, this calibration curve the diffusion coefficients values for ferrocene and ferrocenyl acetic acid could be determined based on their  $t_{\text{mc}}$  values. However, the case of ferrocene served to illustrate the danger of using too large values for  $t_{\text{pulse}}$ . Indeed, ferrocene reported to have a larger diffusion coefficient,  $2.24(\pm 0.07) \times 10^{-5} \text{ cm}^2/\text{s}$ (28), than those used to construct the calibration line in Figure 12. Using the same pulse width  $t_{\text{pulse}} = 61$  ms gave a  $t_{\text{mc}}$  value of 0.93 s, which the calibration curve translated as diffusion coefficient of  $1.22(\pm 0.03) \times 10^{-5} \text{ cm}^2/\text{s}$  for ferrocenium, being then in strong contrast with the value of  $2.24(\pm 0.07) \times 10^{-5} \text{ cm}^2/\text{s}$  reported in the literature (38). A similarly wrong outcome was obtained for the diffusion coefficient of ferrocenyl acetic acid. Using again  $t_{\text{pulse}} = 61$  ms afforded  $D = 1.10(\pm 0.05) \times 10^{-5} \text{ cm}^2/\text{s}$  upon using calibration line in Figure 12 while a diffusion coefficient of  $1.92 \times 10^{-5} \text{ cm}^2/\text{s}$  is predicted for ferrocenyl acetic acid based on its molecular weight(39). In fact, the two species diffuse too rapidly for using a  $t_{\text{pulse}} = 61$  ms value in the ETOF experiments.  $t_{\text{pulse}}$  was thus decreased from 61 ms to 24 ms in order to satisfy the above criterion ( $t_{\text{pulse}}/t_{\text{mc}} = 0.037$ ). Using  $t_{\text{pulse}} = 24$  ms afforded a diffusion coefficient value of  $2.4(\pm 0.1) \times 10^{-5} \text{ cm}^2/\text{s}$  (see e.g. the relative positions of the two corresponding data point for ferrocene in Figure 12) for ferricinium being within 7% of the reported literature value for that of ferrocene at room temperature in acetonitrile(28). This illustrates the importance of verifying a posteriori that the  $t_{\text{pulse}}/t_{\text{mc}}$  criterion was satisfied after  $t_{\text{mc}}$  has been determined for the selected  $t_{\text{pulse}}$  value.

### **Aqueous Diffusion Coefficient Calibration Curves Can Extend to Organic Solvents**

The calibration curve in Figure 12 was constructed using literature data values ( $D$ s) and measured ones ( $t_{\text{mc}}$ s) in 0.1 M KCl. However, the value of the constant  $K$  is an intrinsic characteristic of the micro-device and not of the redox system investigated or the electrolyte medium used. Hence, the calibration curve in Figure 12 is expected to be valid in all electrolytes even if it was constructed using specific redox systems in 0.1M KCl. To test the validity of this prediction, ETOF measurements were performed in 0.1 M TBAPF<sub>6</sub> / acetonitrile for ferrocene, Ru(bpy<sub>2</sub>)Cl<sub>2</sub> and the VO<sub>2</sub>(dipic) dimer and their diffusion coefficients compared

to those reported in the literature or determined using RDE for the VO<sub>2</sub>(dipic) dimer. The corresponding results are reported in Figure 13 in which the regression line is that determined in 0.1 M KCl in Figure 12.



**Figure 13:** Use of the calibration curve reported in Figure 3 for 0.1 M KCl to determine diffusion coefficients from tmc values determined in acetonitrile with 0.1 M TBAPF6 as the supporting electrolyte for Ferrocene, Ru(bpy)<sub>2</sub>Cl<sub>2</sub>, VO(acac)<sub>3</sub>, and VO<sub>2</sub>(dipic); for Ferrocene, t<sub>pulse</sub> = 24 ms was used while t<sub>pulse</sub> = 61 ms was used for the three other species. Blue circles represent the set of data determined for Ru(NH<sub>3</sub>)<sub>6</sub><sup>3+</sup>, Fe(CN)<sub>6</sub><sup>3-</sup> and Fe(CN)<sub>6</sub><sup>4-</sup> in 0.1 M KCl aqueous solution, while orange circles represent the set of data determined in in acetonitrile / 0.1 M TBAPF6.

As indicated above, the ETOF result for ferrocene matched perfectly the literature data(38). The same was true for Ru(bpy)<sub>2</sub>Cl<sub>2</sub>(39),  $D = 9.3(\pm 0.4) \times 10^{-6} \text{ cm}^2/\text{s}$ . This validates our claim that calibration curves constructed using standard mediators in aqueous electrolyte, can be used to determine diffusion coefficients in other electrolytes including organic media.

For the VO<sub>2</sub>(dipic) dimer a  $3.9(\pm 0.2) \times 10^{-6} \text{ cm}^2/\text{s}$  value was obtained based on ETOF (see Figure 13) being in excellent agreement with that,  $3.94 \times 10^{-6} \text{ cm}^2/\text{s}$ , determined using the slope of the Koutecky-Levich plots (equation 1, with a 95% confidence level) deduced from RDE experiments performed on the

VO<sub>2</sub>(dipic) dimer. This latter example evidenced that the ETOF method is not at all restricted to A/B redox systems that involve a single reversible electron transfer step but is applicable to more complex kinetic sequences (see also below the case of VO(acac)<sub>2</sub>).

### **Application to VO(acac)<sub>2</sub> in acetonitrile / 0.1 M TBAPF<sub>6</sub>**

As recalled in the Experimental section, the electrochemical oxidation of vanadyl acetylacetonate, VO(acac)<sub>2</sub>, samples is reported to easily lead to VO(acac)<sub>3</sub> (30). This was confirmed through cyclic voltammetry since the VO(acac)<sub>2</sub> oxidation wave was irreversible and that of VO(acac)<sub>3</sub> alone was observed during the backward CV scan.

The present ETOF method amounts to determine the diffusion coefficient of the species generated at the generator and not of that present in the solution bulk. So, the VO(acac)<sub>2</sub>/VO(acac)<sub>3</sub> dichotomy offered a good opportunity to test the applicability of ETOF when the electron stoichiometry of the electrochemical reaction producing the species of interest (here VO(acac)<sub>3</sub>) is complex and a priori unknown under analytical conditions.

The electron stoichiometry of VO(acac)<sub>2</sub> oxidation at the generator electrode is expected to generate the cation VO(acac)<sub>2</sub><sup>+</sup> as the primary intermediate (the oxidation numbers of the vanadium center are noted in the following to help checking the father-son sequences):



The highly chemical instability of this primary intermediate was confirmed by the chemically irreversible CV oxidation wave of V<sup>IV</sup>O(acac)<sub>2</sub>. Owing to the reported instability of V<sup>IV</sup>O(acac)<sub>2</sub> in the presence of traces of its oxidation product V<sup>V</sup>O(acac)<sub>3</sub> (30) and the irreversibility of the CV wave of V<sup>IV</sup>O(acac)<sub>2</sub>, V<sup>V</sup>O(acac)<sub>2</sub><sup>+</sup> is expected to undergo a rapid father-son follow-up mechanism(18):



Equation 9 affords the neutral V<sup>V</sup>O(acac)<sub>3</sub> species whose diffusion over the generator-collector(s) gap would then be measured instead of that of the primary oxidation intermediate VO(acac)<sub>2</sub><sup>+</sup>. Note that the

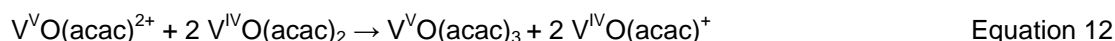
overall electron stoichiometry of the mechanism in equations 8,9 is equal to 0.5 since one unreacted  $V^{IV}O(acac)_2$  is consumed by the ligand exchange in equation 9:



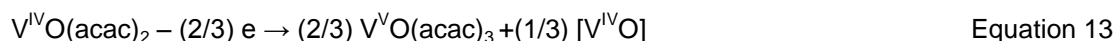
This global stoichiometry would even be more complex if  $V^{IV}O(acac)^+$  is oxidizable at the oxidation wave of  $V^{IV}O(acac)_2$ :



as suggested by the absence of a reduction wave that could be ascribed to this species in the CV of the  $V^{IV}O(acac)_2 / V^VO(acac)_3$  redox system. This may look unlikely at first glance owing to the coulombic repulsion due to its positive charge. However, one expects that the resulting  $V^VO(acac)^{2+}$  species reacts much faster with the parent  $V^{IV}O(acac)_2$  species than does  $V^VO(acac)_2^+$ :



Equation 12 may thus displace the oxidation wave of  $V^{IV}O(acac)^+$  much before its standard potential(41) so that it may occur before that of  $V^{IV}O(acac)_2$ , so that a global father-son ECECE sequence takes place at the potential of the  $V^{IV}O(acac)_2$  oxidation wave, thus leading to the global stoichiometric outcome in equation 13:



where  $[V^{IV}O]$  represents unreacted ' $V^{IV}O$ ' vanadyl species and electrochemically inert over the -2V, 2 V potential range voltammetrically explored.

In other words, depending on the extent of the conversion in equation 13, the electron stoichiometry of  $V^{IV}O(acac)_2$  is expected to range between 0.5 and 2/3 of an electron per vanadyl center. This evidences that relying on RDE measurements of  $V^{IV}O(acac)_2$  oxidation and Levich's current limit to extract a diffusion coefficient is necessarily biased by an unknown factor. Indeed, owing to equation 1, the Levich current slope depends on the value of  $p = D^{2/3}n$ , hence the exact  $D$  value is extremely sensitive on the actual value of  $(D = (p/n)^{1.5})$ .

Therefore, this redox system precisely served our purpose of illustrating a great specific advantage of the ETOF method vs. the traditional RDE or SECM approaches. Indeed, using the limiting plateau current in RDE would require a foreknowledge of the value,  $n$ , of the apparent number of electrons exchanged per  $\text{VO}(\text{acac})_2$ . Similarly, current feedback measurements in SECM are highly dependent on the ability of  $\text{VO}(\text{acac})_2$  to be quantitatively regenerated by reduction of  $\text{VO}(\text{acac})_3$  at the collector. Conversely, ETOF measurements rely exclusively on diffusional time durations of the  $\text{VO}(\text{acac})_3$  species and not at all on current intensities.

The diffusion coefficient of  $\text{VO}(\text{acac})_3$  in acetonitrile with 0.1 M  $\text{TBAPF}_6$  as the supporting electrolyte was thus deduced experimentally based on the corresponding  $t_{\text{mc}}$  value using the calibration curve without any hypothesis about its exact nature. This afforded a value of  $1.71(\pm 0.08) \times 10^{-6} \text{ cm}^2/\text{s}$  (Figure 4). In agreement with the complex mechanism detailed above, such value is in total disagreement with the estimated diffusion coefficient based on the molecular weight of  $\text{V}^{\text{IV}}\text{O}(\text{acac})_2$ , viz., of its primary oxidation intermediate  $\text{V}^{\text{VO}}\text{O}(\text{acac})_2^+$ . Conversely, a diffusion coefficient has been reported for  $\text{V}^{\text{VO}}\text{O}(\text{acac})_3$  (3), viz.,  $2.3 \times 10^{-6} \text{ cm}^2/\text{s}$ , being only 25% larger than the value measured by ETOF. Hence, it is presumable that the species whose diffusion coefficient was determined by ETOF was  $\text{V}^{\text{VO}}\text{O}(\text{acac})_3$ .

## Conclusions

This paper shows that ETOF is an elegant, effective, and simple method that can be used to easily determine the diffusion coefficient of the product of a redox reaction even when this is formed through a complex mechanistic sequence. Indeed, the ETOF approach relies only on the duration of the diffusional transport between a generator electrode where the species of interest is formed and one or two collector electrode(s) at which it can be detected electrochemically and not at all onto electrochemical current intensities. This contrasts with other approaches such as that using the Levich limit of RDE currents or as SECM that relies on current feedback magnitude. This unique advantage of the ETOF approach has been successfully illustrated in the context of vanadyl acetylacetonate,  $\text{VO}(\text{acac})_2$ , oxidation that underwent a complex mechanism with unknown fractional electron stoichiometry. Irrespective of this value, the diffusion coefficient of  $\text{VO}(\text{acac})_3$ , the stable oxidation product of this

complex oxidation wave, could be determined and shown to be in good agreement with a previously reported value.

A second unique advantage of the ETOF method is that the relationship between the time of flight,  $t_{mc}$ , measurements and the sought diffusion coefficients,  $D$ , is independent of the electrolyte / solvent media in which the ETOF measurements are performed, but depends exclusively on the geometrical characteristics of the micro-device used for ETOF experiments. This has been illustrated by establishing a calibration curve,  $\sqrt{t_{mc}} \propto 1/\sqrt{D}$  (Figure 3), based on the measurements of  $t_{mc}$  values in 0.1 M KCl for a series of species whose diffusion coefficients were reported in the literature, and validating its use in 0.1 M TBAPF<sub>6</sub> / acetonitrile (Figure 13).

Albeit these important and unique advantages, it is emphasized that a proper use of the ETOF method requires that the duration,  $t_{pulse}$ , of the potential pulse applied to the generator electrode is not excessively long compared to the experimental time of flight,  $t_{mc}$ . This is an essential criterion for guaranteeing a total absence of diffusion cross-talk between the collector electrode(s) and the generator one which, otherwise, would introduce significant biases in the ETOF measurements based on equations 2 and 3. This important caveat has been illustrated based on two ETOF measurements performed for ferrocene in 0.1 M KCl at two different  $t_{pulse}$  values, one of which was too large to satisfy this criterion and provided a too small  $D$  value compared to published ones.

In this respect it is worth to mention that the whole ETOF experiments can be readily modeled using commercial softwares such as Comsol so as to validate quantitatively the experimental calibration curve corresponding to a given micro-device, as well as for establishing quantitatively the criterion fixing the upper limit of  $t_{pulse}$  values for a given experimental value of the time-of-flight  $t_{mc}$ .

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## **Future Prospects and Conclusions**

### **Recalibration of Oxygen Sensing Electrodes**

The development of a tool for determining performance changes in both oxygen and enzyme electrodes would make a significant impact on the development of sensors that are constructed with membrane coverings. Since so much of the performance of these sensors depends upon membrane permeability, a new method to determine the diffusion rate of molecules through membranes would have a broad impact. These advancements would lay the ground work for obtaining more accurate results in clinical applications involving implanted electrodes, providing a novel way to recalibrate electrodes in vivo. This should allow for more accurate real-time measurements of analytes in clinical studies. It could also be used in other situations where the electrodes are damaged by the in situ situation, such as probing for oxygen in deep water or cooling water in a power plant.

Electrodes have been used to determine the oxygen concentration in vivo. The most common electrodes used for this are based on the Clark type electrode. At a fundamental level, the sensitivity of the Clark type electrode is set by the permeability of the membrane that covers the electrode(1). In addition, when used in vivo, the membrane layer becomes fouled by non-specific adsorption, and the immune response will eventually encapsulate the electrode. It can take a week for encapsulation to occur and before that changes in membrane permeability with protein fouling cause changes in sensitivity, not only for the Clark type oxygen electrode but for many types of electrochemical sensors that use a membrane coating(1).

The present method for calibration of electrodes in vivo and in situ electrodes is to use an external calibration, then implant the electrodes into the desired environment. The calibration used in vitro is expected to hold when used in vivo. This is rarely the case because the matrix of the analyte is often dissimilar to the matrix of the calibration solutions(2). When placed in living tissue, non-specific adsorption onto the membrane changes the permeability of the membrane and the rate at which analytes can approach the electrode's surface, thus shifting in the sensitivity. Not correcting for the shift in sensitivity after implantation causes errors in the concentrations reported by the probe(1). Quantitatively measuring changes in membrane permeability could be used to improve the accuracy of the calibration of electrodes that are being used for in vivo analysis.

There are non-electrochemical ways to accomplish this, such as light scattering spectroscopy, NMR spectroscopy, and by using radioactive tracers(3). However, diffusion through membranes is usually determined using electrochemical methods, the most popular being cyclic voltammetry at rotating disk electrodes(RDE's)(4). Impedance spectroscopy(5), non-rotated cyclic voltammetry(6), and wall jet chronoamperometry(7) are also popular choices.

### **A Brief Comparison of ETOF to Other Methods for determination of Diffusion through membranes**

Electrochemical time of flight (ETOF) is an electrochemical technique in which an analyte is generated either oxidatively or reductively at one electrode, called the generator; it is detected by re-reducing/re-oxidizing the analyte at a second electrode, called the collector. The time,  $t_{\max}$ , that it takes the analyte to travel from the center of the generator to the nearest edge of the collector,  $d$ , is the time of flight. By measuring the time of flight, one can then determine the diffusion coefficient through the bulk solution. ETOF could also be developed to determine the diffusion coefficient of analytes through a membrane that coats the surface of the electrode. Quantitating diffusion coefficients of analytes through membranes would be useful for investigating changes in membrane permeability after biofouling. It would provide a method of in situ recalibration of oxygen electrodes, at least for oxygen as previously reported with macro electrodes(1). Recalibration of an oxygen sensor has been accomplished using macro electrodes, but not with ETOF. However presently there are gaps in the literature that have previously prevented one from being able to utilize ETOF to its fullest extent and in this dissertation I presented my work that turns ETOF from a novelty into a full-fledged technique that should enable it to become a standard method for the determination of diffusion coefficients.

However one place where ETOF has not been previously used is in the determination of diffusion coefficients and permeabilities through membranes that have been polymerized on the surface of an electrode. Determining the permeability and diffusion of electroactive molecules through membranes has been done with cyclic voltammetry (CV), rotating disk electrodes (RDEs)(8-14) or rotating ring disk electrodes (RRDEs)(15). Some qualitative work can be done by just looking at a change in the CV (16-18) of a probe molecule. More quantitative investigations can be done using wall jet chronoamperometry(7) or impedance spectroscopy(5, 19-23). But the majority of investigations into the

interactions of electroactive molecules and membranes have been done with RDE and/or cyclic voltammetry.

Cyclic voltammetry even without using rotating disk electrodes has been used to evaluate the differences in diffusion through defect sites on a polymer and the diffusion through the rest of the polymer layer(8).

One paper has shown that there are differences in the voltamograms of a probe molecule between through-film transport and pinhole (pore) diffusion(6). Through film transport means that the molecule is dissolved into the film and then diffuses to the electrode, while in pinhole/pore diffusion the molecule transports through evenly distributed pores or pinholes in the polymer across the surface of the electrode. It is very important to know which kind of transport across the membrane is occurring if one wants to be able to determine whether or not a complete coating over the surface of the electrode exists or if there are defects spread across the surface. There are three different types of cyclic voltamograms observed on membrane covered electrodes, with varying scan rate. These responses are a) Randles-Ševčík, b) steady-state and, c) thin layer. Randles-Ševčík responses are characterized by the fact that the calculated peak current density ( $j_{peak}$ ) is a function of square root of the dimensionless scan-rate( $\sigma$ ),  $\gamma$  the diffusion coefficient ratio ( $\gamma = \frac{D_{film}}{D_{bulk}}$ ) of the analyte,  $K_{eq}$  is the equilibrium constant between the analyte in the solution and in the film, Equation 1.

$$j_{peak} = 0.4463K_{eq}\sqrt{\sigma\gamma}. \quad \text{Equation 1}$$

Steady State is the region where peak current as a function of scan rate is described by Equation 2, and the thin film region's peak current is described by the following equation Equation 3.(6)

$$j_{ss} = K_{eq}\gamma \quad \text{Equation 2}$$

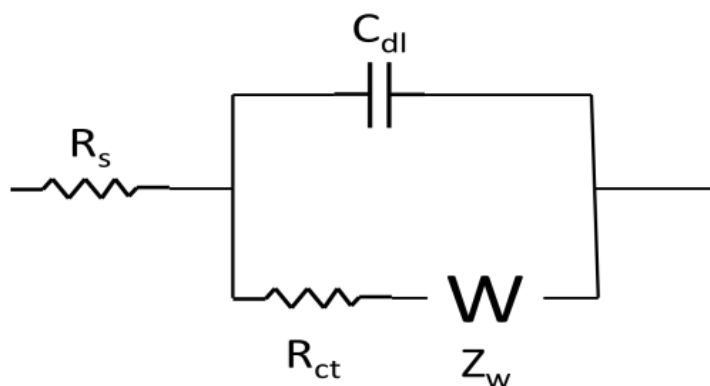
$$: j_{peak} = K_{eq} \left( \frac{\sigma}{4} \right) \quad \text{Equation 3}$$

By watching how the peak current changes as scan rates are shifted over at least three orders of magnitude, one can infer the structure of the electroinactive polymer layer. This is because the thin layer behavior is only exhibited in the case of through-film transport(6). Cyclic voltammetry by itself can also be used for the gathering of preliminary data and qualitative investigations of diffusion(7, 16-18). This can be

done by noting qualitative changes in the shape and height of the peaks of the  $i$ - $E$  curve with membrane thickness or with the presence of a membrane compared to a bare electrode.(17). When polymers are placed over electrodes raised off the substrate (as in the case of MEAs prepared by lithography), CV can reveal slight differences in diffusion through the polymer located at the electrode's edge(16). It can also be used to examine a membranes exclusion to a given probe molecule and how the membrane affects molecules with different polarities and charges, as well as differences in molecular construction of the membrane in question(7).

Data from RDEs can be used to determine diffusion coefficients in bulk solution by using the Levich equation. The Levich equation can also be applied to the diffusion of molecules through membranes that have been polymerized on the surface of the RDE electrode(4). By using a probe molecule, the most common being ferricyanide or hexamine ruthenium, one can investigate the permeability of the membrane to these molecules, either by using a Levich plot or an inverse Levich plot (Koutecky-Levich plot). These methods have been successful for numerous different membranes polyethylene terephthalate(14), poly(allylamine)/poly(acrylic acid) crosslinkages(9), poly diallyldimethylammonium(11), polystyrenes(10, 24), polypyrrole(12), metalopolymers(25),etc..) and various probe molecules other than

ferri-ferrocyanide and hexamine ruthenium, including iodine(9) and quinones(14) in aprotic solvents.



**Figure 14:** Circuit diagram of the Randles Equivalent Circuit,  $R_s$  is the solution resistance,  $C_{dl}$  is the capacitance of the double layer,  $R_{ct}$  is the resistance to charge transfer, and  $Z_w$  is the Warburg element.

Another method to determine the permeability/ diffusion of molecules through a membrane is impedance spectroscopy. A small sinusoidal potential is used, as opposed to the much larger potential window that is used by cyclic voltammetry(5). A Randles circuit is used to model the data on the polymer film that

coats the electrode(5), Figure 14. The resistance to charge transfer circuit element in the Randles' Circuit,

$R_{ct}$  which is directly related to the ability of analyte to transfer an electron, or the accessibility of an underlying electrode(5). The data can be graphed in Nyquist plots, to determine  $R_{ct}$ , which can be used to calculate the conductivity of the polymer ( $\sigma$ ). The diffusion coefficient,  $D$ , can then be calculated using the Nernst-Einstein equation, Equation 4, where  $k$  is the Boltzman constant,  $T$  is the temperature at the measurement.  $[E]$  is the concentration of the electroactive compound, and  $q$  is the ionic charge of the diffusing species(5).

$$D = \frac{\sigma kT}{[E]q^2} \quad \text{Equation 4}$$

The problem with this equation is that it can only be used for compounds that are ions in solution.

Impedance spectroscopy has also been used to look at the impedance of mass transport of probe molecules through polyelectrolyte “brushes” using a different equation which can be solved for the mass transfer coefficient and thusly the diffusion coefficient of the probe molecules(22). Impedance spectroscopy has also showed the ability to investigate the changes in the electroactivity of the film based on solvent effects(19).

Wall jet chronoamperometry, has been used to determine the permeability and diffusion of probe molecules through silicon films that have been built on the surface of an electrode(7). The wall jet served the same purpose as the motion in the rotating disk electrode system in that both are used to mask diffusion from the bulk solution. When using the wall jet electrode system the equation for the diffusion coefficient and the permeability of the probe molecule is Equation 5, where  $d$  is the thickness of the membrane,  $n$  is the number of exchanged electrons,  $F$  is Faraday’s constant,  $P$  is the permeability of the probe molecule through the membrane,  $D_f$  is the diffusion coefficient of the probe molecule through the membrane, and  $C$  is the concentration of the molecule in the bulk solution(7).

$$\frac{1}{i_{lim}} = \frac{d}{nFA(PD_f)C} \quad \text{Equation 5}$$

None of the above methods have the advantages of ETOF in that it can be performed in situ, making it more powerful when applying it to the diffusion of oxygen or analytes through a membrane on an implanted electrode. This means that one could generate real time data about the changes in the

permeability of a membrane that has been fouled by its surroundings either by non-specific adsorption of proteins to the surface of the membrane, or by the destruction of the membrane in other ways by its surroundings.

## Conclusion

This work shows the broad applicability of ETOF beyond what has been previously shown in the literature. By rearranging the equations from the standard form of the ETOF experiment, a diffusional calibration curve can be constructed to determine diffusion coefficients for many different analytes in many different solvent systems. The caveat is that the pulse width must be short enough to prevent Red/Ox cycling, between the members of the electrode array. The modeling collaboration, with Christian Amatore, Catherine Sella and Laurent Thouin, showed that for our geometry  $\tau_{\text{pulse}}$ , the unitless pulse width, should be less than 0.03, and so as long as this condition is met diffusion coefficients can be determined over a wide linear range. The pulse width,  $t_{\text{pulse}}$  can be changed to allow for optimal signal at the detectors. Many diffusion coefficients in aqueous, viscous, and organic solvent systems, were determined. Future work on the diffusion of molecules through membranes, as well as the study of molecules diffusing through ionic liquids, and other novel solvent systems would be possible.

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## Appendix 1: Detailed Device Construction

**Introduction to Labview:** LabVIEW is now a standard interface language developed by National Instruments (Austin, TX) that has been used to control thousands of commercial instruments. LabVIEW is a visual programming language allowing the user to create programs which can control computer interfaces, such as the National Instrument cDAQ-9174, NI Compact DAQ chassis. NI cDAQ-9174 is used to control digital input and output signals (NI 9403 Digital I/O), analog output (NI 9269), and analog input (NI 9215), which all exist as plug-in modules to the chassis. The interface is also equipped with a timer that can be used to control events, such as the application of potential pulses, or a sequence of pulses. As a visual programming language LabVIEW uses glyphs and other representations of programming concepts to allow for a simplified, yet still powerful and versatile programming platform. There is also a large open-source community of canned programs that others have already built that can be either refactored or used as guidelines for building your own program.

There are five main groups of programming concepts used in LabVIEW: loops/structures, subVIs, controls, indicators, and constants. Using LabVIEW, a virtual instrument (VI) is constructed which is composed of several sub VIs that make up the majority of glyphs in the program. Loops/Structures are usually represented as boxes with special borders such that subVI's can be placed into the loop. Controls connect to the user interface which appears as a virtual representation of the front panel of an instrument. These controls can be represented as boxes for user supplied data, dials, slides, etc. Indicators are another way that the program interacts with the user; they usually exist in the form of lights, graphs, or other forms of data recording or indicating objects that appear on the front panel. Constants are a way to program a binary (true, false) or constant values into the program.

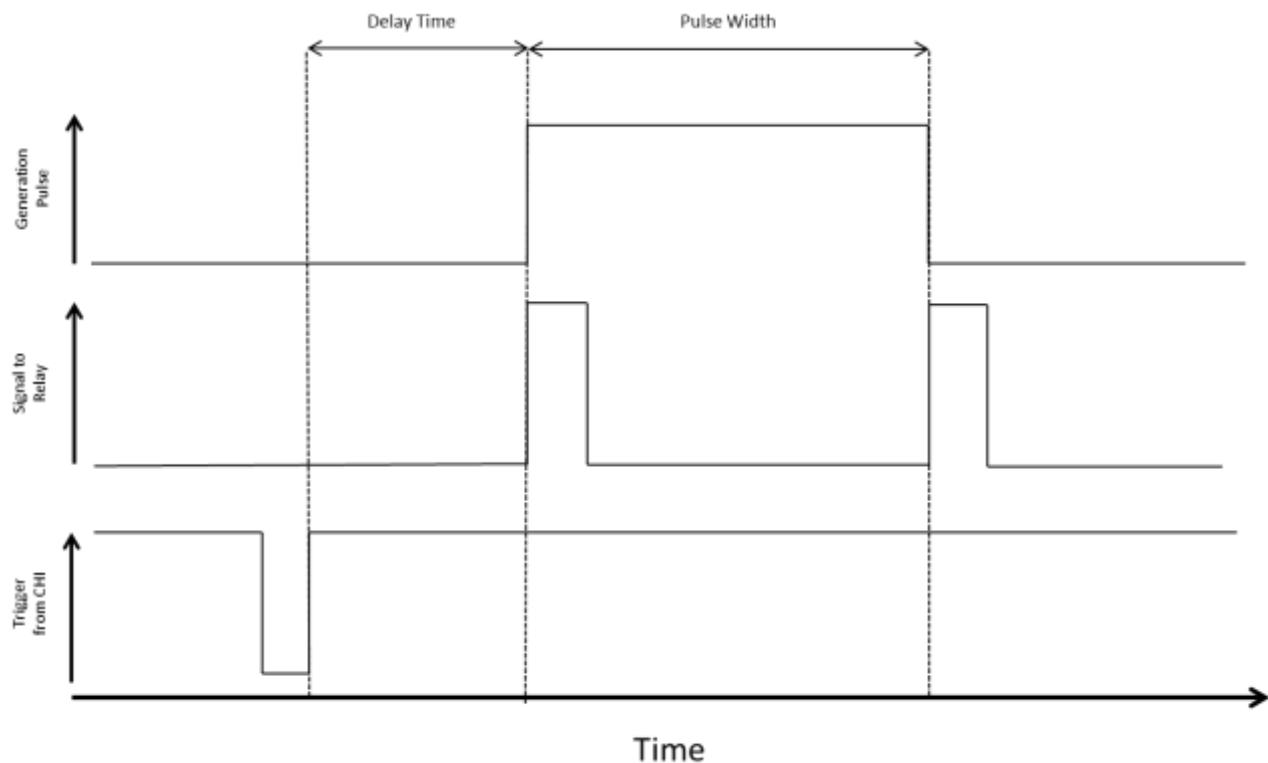
As a dataflow language, LabVIEW allows one to program tasks to take place in parallel instead of sequentially. In fact in most cases one has to go out of their way to program a task so that it takes place in a sequence instead of in parallel. This is because any piece of code will initiate as soon as all of the inputs have arrived. If two programs receive all of their inputs at the same time then both programs will run simultaneously. There are ways to work around this, which is necessary for the task here. One is to wire the errors from one program to the next so that the second program will wait until the first program



has stopped working to do its task because it needs the input from the error cluster to have arrived before it will start working. One could also use the sequence structure which is a brute force way of programming a sequence of actions for LabVIEW to perform by programming each of the individual code pieces in frames, which are then activated from right to left until the program is finished. To control the interface that performs the ETOF experiment, sequential operations were forced upon the inherent parallel structure.

If required, a complete multi-stat control for ETOF experiments would be possible using LabVIEW. The VI software can accommodate any situation that occurs in real time programming and instrument control. A few of these capabilities are outlined in this section. Parallelism can also cause problems when there are codependent applications running because of the fact that each of the applications may need to share data back and forth between one another and this is a job that requires the use of variables, but one must be careful when programming with variables. This is because they break the dataflow paradigm that is used for the majority of LabVIEW programming and this can cause unexpected issues if one is not careful. As such unexpected issue is a race condition, which occurs when a variable is being written by two or more sources in parallel. This is because either source might overwrite the variable before it has to be transferred to the desired location to be used. In any case it is best to have variables only written in one place while they can be read in multiple places in the code. There are also other programming structures that are very useful to use. Suppose one wants the program to change based on values or other types of conditions or states. For this, a state based architecture is used, which incorporates the code and how the code is supposed to be used to change between the states of the device. To repeat a given application until finished, a while loop with a stop condition is used. The stop condition could be a simple Boolean button or it could be code that evaluates until a given condition is met, and then it halts the loop. While loops are useful, they also use a lot of processing time to maintain, taking up a whole processor on a multicore machine just to do iterations. This means that it is wise to program a wait time in between each cycle of the loop so as not to waste the processing power of the computer that the program is running on. Master-slave and producer-consumer are multiple loop design patterns, and in these cases there is a way to transfer the data from one loop to another; via notifiers and queues. Notifiers are used for the master-slave design pattern in that one loop, the master, sends a notifier to the slave loop as to what task the slave loop should be performing. This is especially useful for modifying the user interface while the

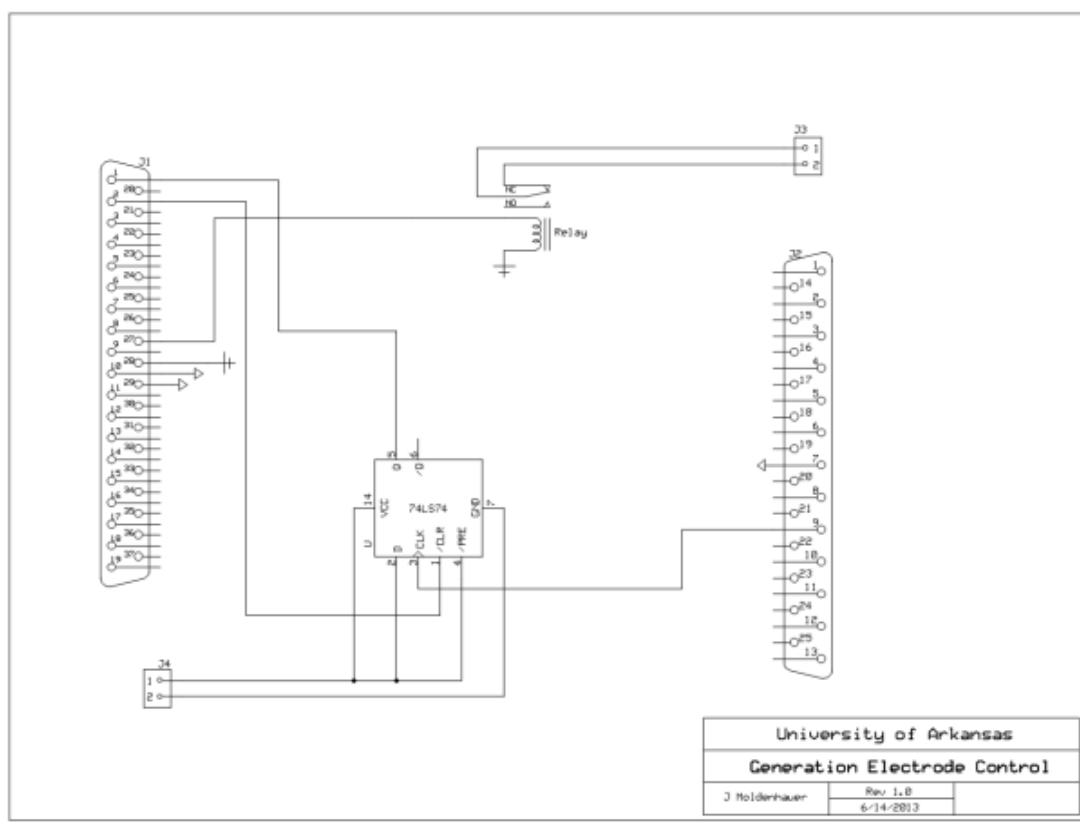
instrument is recording data. This design pattern works well only when the slave loop completes its tasks faster than the master loop will. Otherwise the master loop has the chance of accidentally writing over its own instructions to the slave loop. Conversely producer and consumer design pattern is for dealing with data as it is created, working using queues to transfer the data between each of the loops, buffering the data it receives from the producer loop as it is being transferred to the consumer loop. “Event structures” are useful for programming more complicated state- based-machines because the event structures will automatically respond to events on the user interface. This is a method that most likely works best when one is using a master-slave design pattern. That way the event structure can be in the master loop and has



**Figure 15:** A timing diagram of the operation of the generation electrode controller, showing the outgoing trigger from the CHI 750a and the two pulses from the NI 9402 to open and close the relay.

simplistic programming to respond to the interactions with the user interface while the slave loop is running data collection or data generation. This is one of the caveats and recommendations for using event structures instead of other state based design patterns that are covered at length in the LabVIEW help.

**Initial Development of Program and Device Design:** The LabVIEW environment can be used easily to develop a program that would be triggered by CHI potentiostat and controlling a relay for the generation pulse at the second working electrode. Figure 15, It shows a timing diagram for opening and closing relay in the generation electrode controller. The relay can be triggered to open and close in response to timers that are triggered by the “start scan” pulse from the CHI 750a potentiostat.



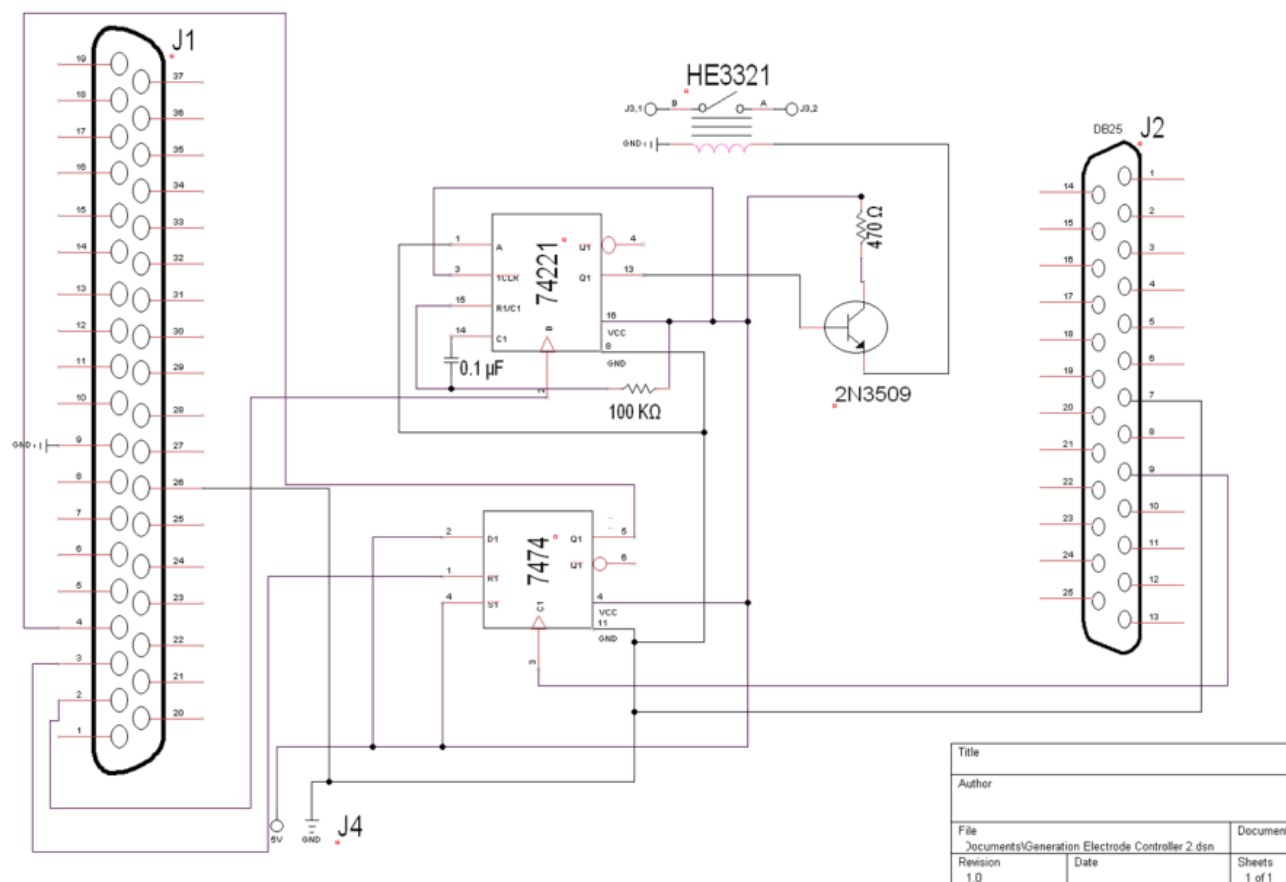
**Figure 16:** The circuit diagram of the GEC illustrating all of the connections made by the DB connectors to the chips on the card. J1 is the connector between the generation electrode controller and the NI digital I/O module. J2 is the connector between the generation electrode controller and the CHI. J4 is the connector between the NI analog output module that was used to power the generation electrode controller.

The start scan trigger is caught by a D-triggered flip-flop, Figure 16, so there is no worry about the program missing the outgoing trigger signal from the CHI. The program written in LabVIEW will use a sequence structure so that each event will occur in sequential order, first resetting the flip-flop, then waiting for the scan start pulse from the CHI. In the first software design the timers control a generation pulse by closing and then re-opening the relay, attached to the secondary electrode lead. One of the analog outputs from the NI 9215 is used to power the flip-flop, and the NI 9403 module of a National

Instruments box was used as the digital I/O receiving the CHI start scan trigger, and controlling the relay. The flip-flop and connections were made using wire-wrap construction and a ribbon cable connecting to both the National Instruments interface and the CHI potentiostat. A circuit diagram of this early version of the controller is in Figure 16.

**Improved Program and Device Design:** The triggering of the relay was the most complicated part of this project. Initially the pulse was being triggered by digital high and low signals sent from the NI 9403 module, to open and close the relay. This arrangement gave unpredictable generator pulse widths that varied from what was intended. The minimum possible generator pulse width based on software timers was 40 ms, which was much too long to be useful for the ETOF experiments. Using the electrode array of dimensions, 4 micrometer width electrodes and 4 micrometers separation, causes Red/Ox cycling. Pulsewidths in the range of 7 ms to 15 ms prevent redox cycling from interfering with the time of flight experiment. These times were determined using equations for minimal and average flight time over this distance found in a paper by Amatore.<sup>(1)</sup> Redox cycling can occur when the generator pulse is long enough to allow for analytes to travel back and forth from the generator to collector multiple times, lengthening the time it takes to reach the time of maximum collection. To address this problem a one-shot metastable trigger was added to the circuit as well as a current booster transistor and a 5V power supply. The pulse widths are exact, and can be even smaller than 7 ms, but timing is now set by a hardware

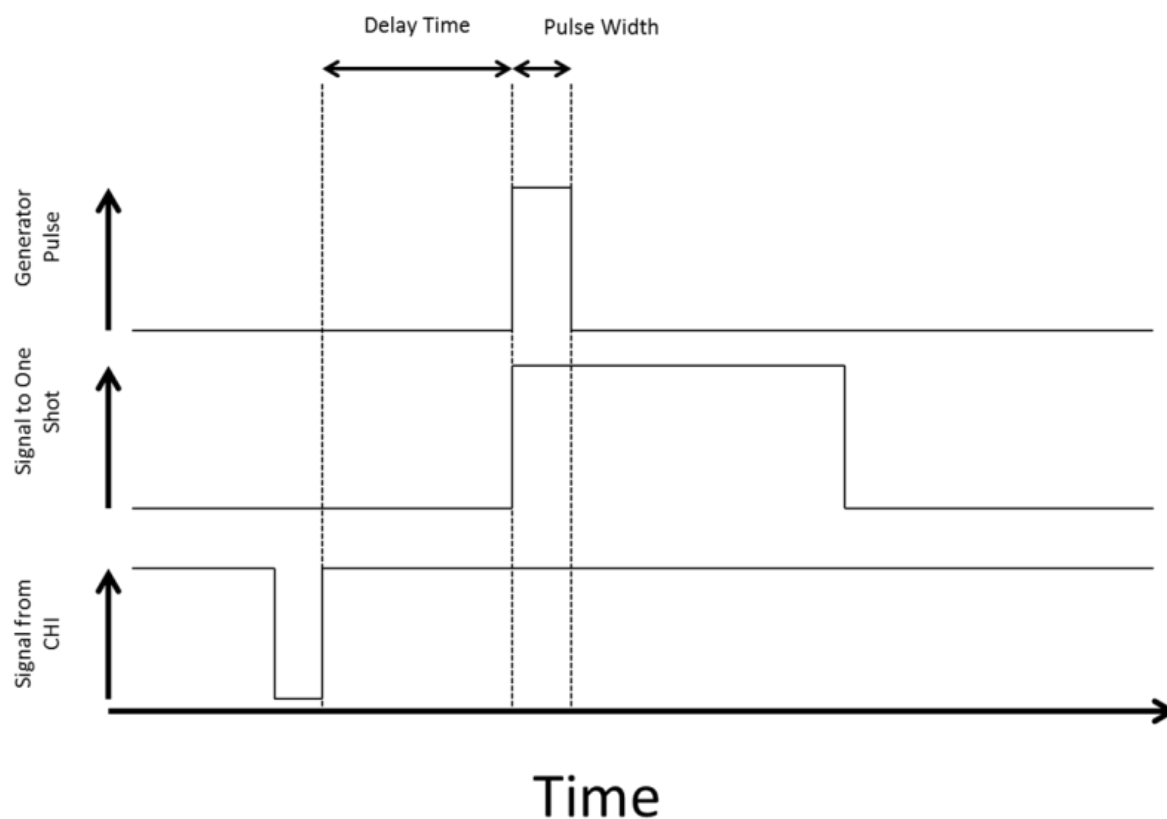
change, selecting a resistor for the correct generator pulse width desired, Figure 17.



**Figure 17:** The improved circuit diagram for the generation electrode controller version 2 includes the placement of the 74221 metastable trigger and the 2N3509 transistor. Junctions remain the same as on the previous design where J1 connects to the NI digital I/O module, J2 connects to the CHI potentiostat, and J4 connects to an external power supply.

In the original design, the power was supplied by the NI 9215 analog output module, but the newer version features a 5V power capable of powering all the chips on the wire-wrap board including the addition of the one-shot metastable trigger. This 5V power source was constructed from an old cellphone charger, to convert the AC outlet to the direct current needed to run the generation electrode controller. Another problem that arose was that the one-shot trigger did not have enough output current to be close the relay, so the relay was briefly replaced with a CMOS switch relay, but it was abandoned because it could not pass negative potentials when the switch was closed. The problem was solved by the insertion of a transistor to boost the current from the one-shot so that there was enough current to close the relay.

These changes also lead to changes in the timing diagram updates now shown in, Figure 18.

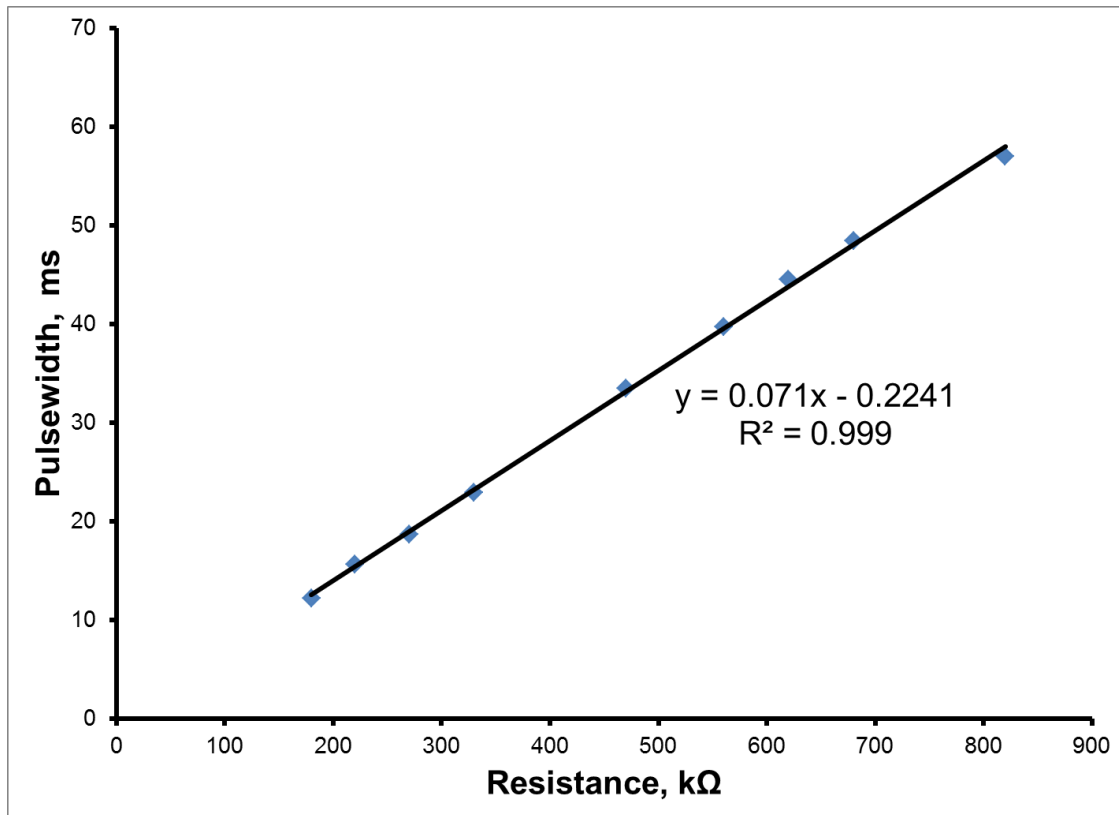


**Figure 18:** The new timing diagram for the Generation Electrode Controller version 2 taking into account the changes based on the use of the one-shot metastable trigger.

In the course of improving the generation electrode controller so that it could utilize the digital I/O from the NI 9403 module, the generation electrode controller needed to be rewired several times, and in the course of that rewiring, ground issues were created and had to be worked out. This led to the wire-wrap being color coded green for digital ground, blue for analog ground, and red for +5 V. The addition of additional grounding and the uncertainty of wire wrap connections called for a printed circuit board. The wire wrap prototype was upgraded to a printed circuit board, but the first iteration lacked the ability to rapidly change resistors to control the pulse width of the generation pulse for ETOF experiments. Because a constant generator pulse width is only applicable over a very small range of diffusion coefficient values, because of redox cycling. So as diffusion coefficients become faster one needs to be able to change out resistors to allow for a broader range of diffusion coefficients to be determined. So in

order to change the generator pulse width the resistors need to be changeable or at least tuneable with a trimpot.

In the absence of a trimmer potentiometer or trimpot, one needs to have a function to relate size of a resistor to the pulse width. This function was generated empirically by measuring the generator pulse width using an oscilloscope and recording the resistor values, to make a linear plot, Figure 19, over a range of resistors that seemed to produce a reasonable range of pulse widths for the electrode array dimensions that were available. As the pulse width increases one needs to increase the gap to maintain linearity. As there are fewer options on how to change the gap of the electrodes on a printed array, ability to change the pulse width resistor was of highest importance.



**Figure 19:** Based on the slope of this line one can determine the pulse width generated by the one shot based on the size of the resistor at a constant capacitor value of 0.1  $\mu\text{F}$ . The resistors chosen provide a range of about 100 ms for pulse widths for the generator electrode.

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